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Higgs self-coupling measurements using deep learning in the $b\bar{b}b\bar{b}$ final state

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ABSTRACT: Measuring the Higgs trilinear self-coupling λ_{hhh} is experimentally demanding but fundamental for understanding the shape of the Higgs potential. We present a comprehensive analysis strategy for the HL-LHC using di-Higgs events in the four b -quark channel ($hh \rightarrow 4b$), extending current methods in several directions. We perform deep learning to suppress the formidable multijet background with dedicated optimisation for BSM λ_{hhh} scenarios. We compare the λ_{hhh} constraining power of events using different multiplicities of large radius jets with a two-prong structure that reconstruct boosted $h \rightarrow b\bar{b}$ decays. We show that current uncertainties in the SM top Yukawa coupling y_t can modify λ_{hhh} constraints by $\sim 20\%$. For SM y_t , we find prospects of $-0.8 < \lambda_{hhh}/\lambda_{hhh}^{\text{SM}} < 6.6$ at 68% CL under simplified assumptions for 3000 fb⁻¹ of HL-LHC data. Our results provide a careful assessment of di-Higgs identification and machine learning techniques for all-hadronic measurements of the Higgs self-coupling and sharpens the requirements for future improvement.

KEYWORDS: Higgs Physics, Beyond Standard Model

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1 Introduction

Discovering Higgs boson pair production $pp \rightarrow hh$ opens the only direct laboratory probe of the Higgs trilinear self-coupling λ_{hhh} , which is a principal goal of the LHC and its upgrades [1–7]. Measuring λ_{hhh} is critical for characterising the dynamics of electroweak symmetry breaking that could be modified by beyond the Standard Model (BSM) physics [8–38]. Recent experimental [39–45] and phenomenological [46–48] advances in the four bottom quark channel $hh \rightarrow 4b$ suggest that this final state is competitive for discovery [49–53]. Current 95% CL limits from ATLAS [54] (CMS [55]) on the dominant gluon fusion cross-section reach 6.9 (22.2) times the SM, with the $4b$ channel being second most constraining in ATLAS.

For direct self-coupling λ_{hhh} constraints, the present best combined 95% CL limit with respect to the SM value $\lambda_{hhh}^{\text{SM}}$ is $-5.0 < \lambda_{hhh}/\lambda_{hhh}^{\text{SM}} < 12.0$, where $4b$ is among the most competitive channels for values near $\lambda_{hhh}^{\text{SM}}$ [54]. ATLAS also reports tighter limits of $-2.3 < \lambda_{hhh}/\lambda_{hhh}^{\text{SM}} < 10.3$ when combining with indirect constraints from single Higgs channels [56]. In the $4b$ channel, these constraints and High Luminosity LHC (HL-LHC) projections [57, 58] only optimise for the SM coupling and without using events with boosted Higgs bosons. However, non-SM λ_{hhh} values can modify Higgs boson kinematics such that these analyses are no longer optimal for constraining λ_{hhh} . Boosted Higgs decays [59–62] are targeted in resonant di-Higgs searches [44, 49], but not for self-coupling constraints. Interestingly, ref. [46] suggests that these boosted reconstruction techniques can improve the discovery significance of $hh \rightarrow 4b$. Recent progress in $h \rightarrow bb$ analyses [63–70] and the large signal statistics from the high branching ratio $\mathcal{B}(h \rightarrow bb) \simeq 58\%$ motivate use of advanced techniques for signal characterisation.

This paper synthesises these separate advances to present a comprehensive assessment of HL-LHC analysis strategies across all kinematic regimes of the $hh \rightarrow 4b$ channel to evaluate and improve λ_{hhh} constraints. We now discuss the specific novelties of this paper.

We show that analyses optimised for discovery sensitivity of hh with SM couplings are suboptimal for constraining λ_{hhh} at the boundaries of projected limits. This is because SM hh production has more boosted Higgs bosons than scenarios where $\lambda_{hhh}/\lambda_{hhh}^{\text{SM}} \sim 5$ due to the relative contributions of the interfering production amplitudes (figure 1). Reconstructing Higgs bosons with lower boosts is more limited by trigger thresholds, but we nonetheless demonstrate that optimising for signals with non-SM λ_{hhh} rather than SM values can improve λ_{hhh} constraints. Furthermore, due to the small signal-to-background ratios, our sensitivity is limited by systematics, whose size we vary to quantify its impact on λ_{hhh} constraints. Reducing the formidable multijet background can mitigate the impact of systematics, whose suppression relies on modern b -tagging algorithms [71–73]. Our study adopts expected improvements of the impact parameter resolution from inner tracker upgrades crucial for b -tagging.

To further enhance sensitivity, we use neural networks [74–76] for deep learning, which are witnessing widespread applications in particle physics [77–87]. We apply these state-of-the-art techniques to reject backgrounds and improve constraining power of λ_{hhh} . We utilise a recently developed framework called SHapley Additive exPlanations (SHAP) [88] to show what physics information the neural network learns.

Additionally, we investigate the impact of experimental effects such as the finite resolution of jet reconstruction and trigger limitations. This presents a more complete picture of the limitations on the sensitivity of the $hh \rightarrow 4b$ channel beyond previous studies. We also study the composition of the background in detail, including small backgrounds (such as single Higgs production) and the impact of a neural network selection on this composition and shape of background distributions. This is important as uncertainties in background composition affect the experimental systematics of current analyses.

Finally, figure 1 shows that only the triangle diagram contains λ_{hhh} while the box diagram contains two factors of the top Yukawa y_t . Current y_t uncertainties are $\sim 20\%$ from $pp \rightarrow t\bar{t}h$ measurements [67, 69], which we account for in our results due to it being

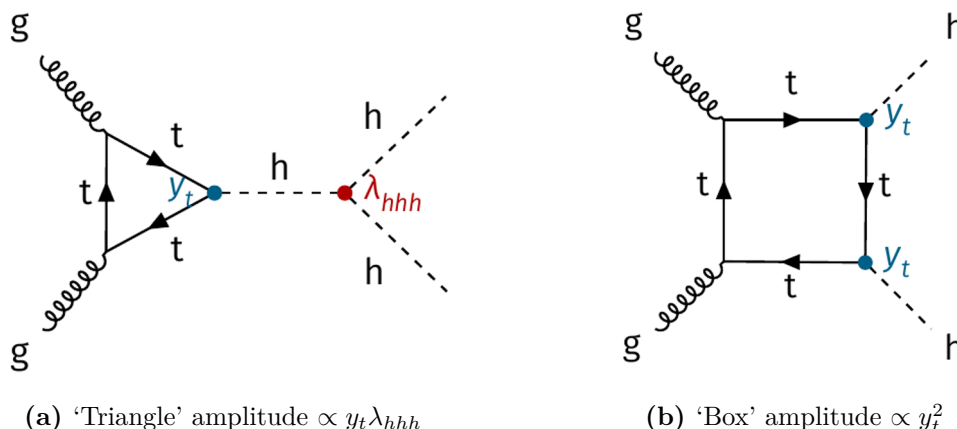


Figure 1. Leading order Feynman diagrams of Higgs boson pair production in gluon fusion. The (a) ‘triangle’ amplitude features the trilinear self-coupling λ_{hhh} and one power of the top Yukawa y_t , and (b) ‘box’ amplitude does not feature λ_{hhh} and has two powers of y_t . These two diagrams interfere destructively in the SM.

sufficiently large to impact λ_{hhh} constraints as demonstrated in ref. [56]. Our y_t -dependent HL-LHC projections of λ_{hhh} constraints allowing equal-footing comparison of sensitivity for different boosted Higgs topologies in the $4b$ channel are new.

While we focus on $4b$ final states, techniques from these studies are readily applied to other channels, such as $bb\tau\tau$ and $bb\gamma\gamma$ that involve $h \rightarrow bb$ decays. Complementary production and decay modes for hh production are explored in refs. [89–107].

This paper is structured as follows. Section 2 outlines the Monte Carlo simulation of the signal and background processes together with discussion of detector emulation. Section 3 describes the baseline and neural network analysis strategies. Section 4 presents the results and statistical analysis before section 5 summarises our conclusions.

2 Signal and background modelling

In the SM after electroweak symmetry breaking, the Higgs self-coupling is determined by the Higgs potential

$$V(h) = m_h^2 h^2 + \lambda_{hhh} v h^3 + \lambda_{hhhh} h^4. \quad (2.1)$$

This paper focuses on measuring the trilinear coupling λ_{hhh} . This will directly test the SM prediction $m_h^2 = \lambda_{hhh} v^2$, and any deviation from this value is indicative of BSM physics. The Higgs boson mass $m_h = 125$ GeV and electroweak vacuum expectation value $v = 246$ GeV are fixed to their independently measured values. Directly probing the quartic self-coupling λ_{hhhh} requires triple Higgs boson production [108–111], which is beyond the scope of this work due to its highly suppressed cross-section. This section details the Monte Carlo simulation of the signal and background used for our analyses. We outline the phenomenology of Higgs pair production along with its simulation (subsection 2.1), background processes (subsection 2.2), and detector emulation (subsection 2.3).

2.1 Higgs pair production

At leading order, the Higgs pair production cross-section has contributions from amplitudes we refer to as ‘triangle’ and ‘box’, shown in figure 1, and their interference. These three contributions scale with the top Yukawa and trilinear self-coupling as

$$\sigma_{\text{triangle}} \sim \lambda_{hhh}^2 y_t^2, \quad \sigma_{\text{box}} \sim y_t^4, \quad \sigma_{\text{interference}} \sim -\lambda_{hhh} y_t^3. \quad (2.2)$$

Sensitivity to λ_{hhh} therefore depends on probing the triangle and interference amplitudes. The comparatively stronger dependence of the total cross-section on y_t highlights the importance of top Yukawa constraints for λ_{hhh} determination. The interference term probes the sign of λ_{hhh} , as it has a relative negative sign due to an additional fermion propagator in the box loop with respect to the triangle. This implies destructive interference for $\lambda_{hhh} > 0$. In the $|\lambda_{hhh}|/y_t \rightarrow \infty$ limit, the total cross-section asymptotes to the same value for either sign of λ_{hhh} as the triangle contribution dominates over the interference.

To modify couplings (λ_{hhh}, y_t) away from SM values $(\lambda_{hhh}^{\text{SM}}, y_t^{\text{SM}})$, keeping all other SM parameters fixed, we use the HEAVYHIGGSTHDM model [114]. This adopts the ‘improved effective field theory’ prescription, which uses gluon-Higgs vertices in the large quark mass limit and is improved by accounting for finite top mass effects as discussed in ref. [115]. Such corrections are important given loop momenta are comparable to the top mass pole. We set couplings to all BSM particles to zero. We vary the top Yukawa y_t while fixing the top mass to $m_t = 172$ GeV. We define $\kappa_\lambda = \lambda_{hhh}/\lambda_{hhh}^{\text{SM}}$ and $\kappa_t = y_t/y_t^{\text{SM}}$ to parameterise variations from the SM couplings, as is conventional in the so-called ‘kappa framework’ [2]. This is a standard prescription, but to consistently study global constraints with other measurements, a full effective field theory treatment is recommended [116–118]; this is deferred to future work.

To sample points in the two-dimensional signal parameter space (λ_{hhh}, y_t) , we employ MADGRAPH 2.6.2 [112, 113]. We consider only the dominant production mode via gluon fusion $gg \rightarrow hh$, inclusive of all Higgs boson decays. The NNPDF3.0 parton distribution functions [119] at next-to-leading order (NLO) are used from the LHAPDF package [120]. We generate 100k Monte Carlo (MC) events per point and calculate cross-sections at leading order (LO) in the strong coupling constant α_s . Figure 2 illustrates the sampled points and LO cross-sections calculated by MADGRAPH, where the SM value is around $\sigma_{\text{LO}}^{\text{SM}} = 16$ fb at $\sqrt{s} = 14$ TeV, and table 1 shows LO cross-sections for example points. For $\kappa_t = 1$, the LO cross-section falls to a minimum of 6.2 fb at around $\kappa_\lambda = 2.5$ due to destructive interference. The cross-section gradient is also shallow around this minimum, rising by only $\sim 15\%$ to 7.0 fb and 7.3 fb for $\kappa_\lambda = 2$ and $\kappa_\lambda = 3$ respectively. This makes λ_{hhh} constraints around such values challenging using inclusive cross-section measurements alone. For the training of the neural network (subsection 3.2), we use an identical generator configuration to produce a dedicated set of high statistics samples with 250k events per κ_λ variation for fixed $\kappa_t = 1$ and use MADGRAPH to decay both Higgs bosons via $h \rightarrow bb$.

We normalise the signal rate to NNLO.¹ We start from the LO MADGRAPH cross-sections and improve their accuracy by applying λ_{hhh} -dependent LO-to-NLO k-factors

¹The cross-section for $\kappa_\lambda = 1$ has recently been calculated to N³LO accuracy [121, 122].

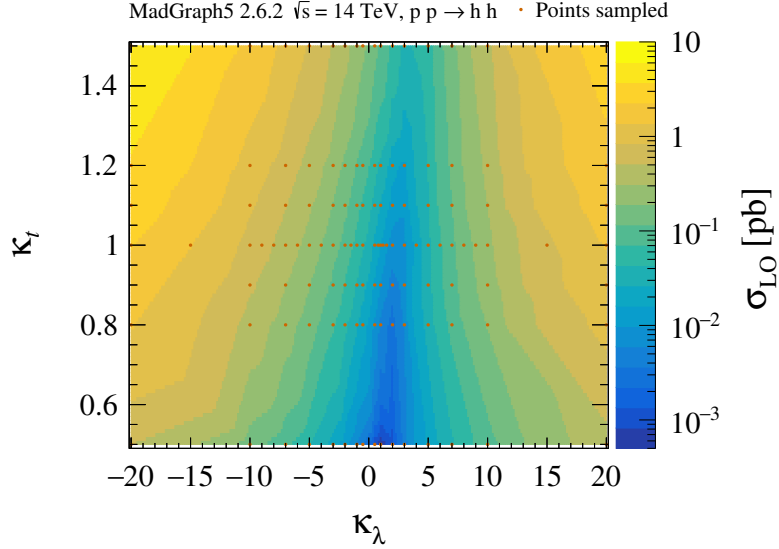
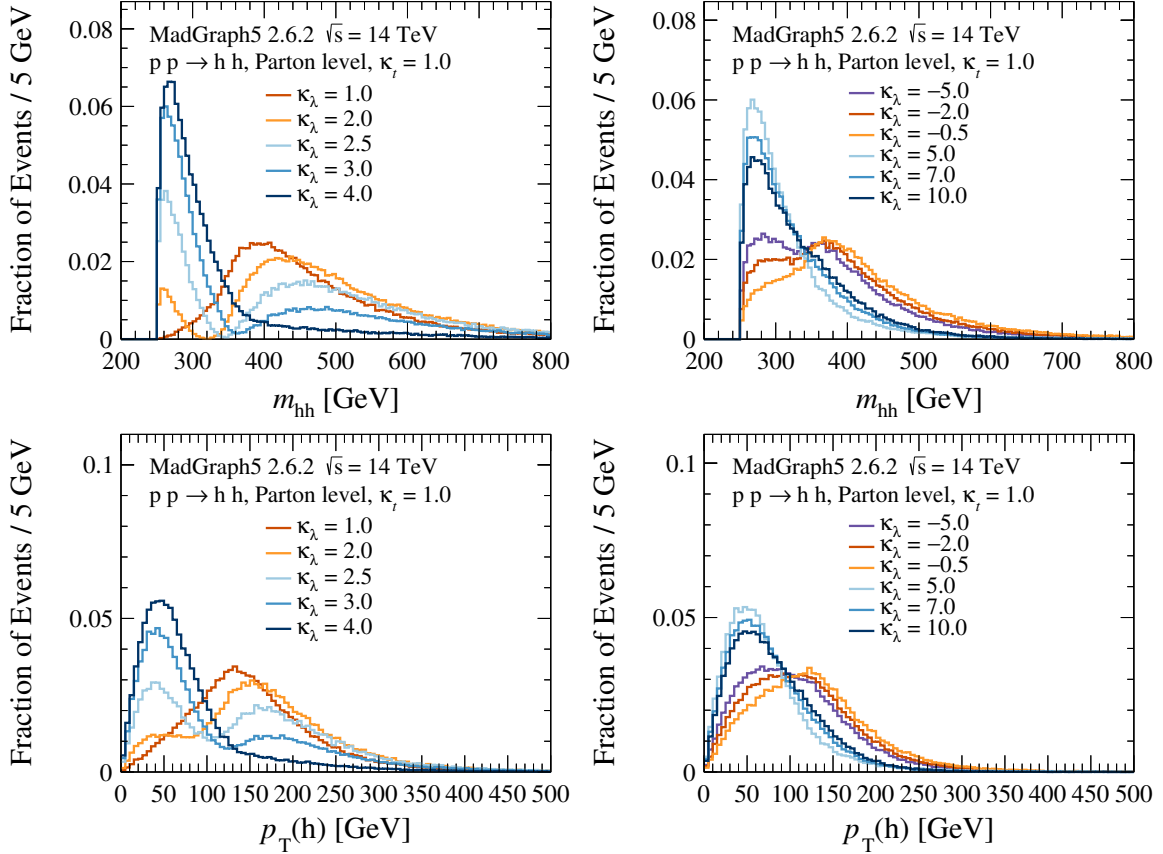


Figure 2. Leading order cross-sections σ_{LO} from MADGRAPH [112, 113] as a function of the variation of Higgs trilinear self-coupling κ_λ and top Yukawa coupling κ_t from their SM values $\lambda_{hhh}^{\text{SM}}, y_t^{\text{SM}}$ using the HEAVYHIGGSTHDM model [114]. The orange markers indicate the points sampled from this parameter space.

$K'_\lambda = \sigma_{\text{NLO}}/\sigma_{\text{LO}}$ and a λ_{hhh} -independent NLO-to-NNLO k-factor $K'' = \sigma_{\text{NNLO}}/\sigma_{\text{NLO}}$ such that $\sigma_{\text{NNLO}} = K'_\lambda K'' \sigma_{\text{LO}}^{\text{MG}}$. For K'_λ , we use the NLO cross-sections for non-SM variations of λ_{hhh} from POWHEG-BOX-v2 [123]. The λ_{hhh} -dependent LO-to-NLO k-factor $K'_\lambda = \sigma_{\text{NLO}}/\sigma_{\text{LO}}$ is 1.66 for $\kappa_\lambda = 1$ and varies by 35% between $-1 < \kappa_\lambda < 5$, ranging from $K'_\lambda = 1.56$ at $\kappa_\lambda = 2$ to 2.14 for $\kappa_\lambda = 5$ [123]. For $\kappa_\lambda < -1$ and > 5 , we smoothly extrapolate the k-factors in ref. [123] to asymptote at 1.95 and 2.1 respectively. In principle, there is a y_t -dependence on these NLO k-factors, but for simplicity we take these to be κ_t -independent. For the NLO-to-NNLO k-factor K'' , the cross-section for $\kappa_\lambda = 1$ at NNLO and next-to-next-to-leading logarithm (NNLL) accuracy [124, 125], with finite top mass corrections to NLO accuracy [126–130], gives $\sigma_{\text{NNLO}} = 39.6 \text{ fb}$ [2]. Based on this, $\sigma_{\text{LO}}^{\text{MG}} = 16 \text{ fb}$ and $K'_\lambda = 1.66$, we derive $K'' = (\sigma_{\text{NNLO}}/\sigma_{\text{LO}}^{\text{MG}})/K'_\lambda = 1.45$. Table 1 shows the LO cross-sections and overall k-factors $K'_\lambda K''$ applied for some example signals.

Turning to the differential distributions, variations in (λ_{hhh}, y_t) away from SM values induce striking features illustrated in figure 3. These are shown at parton level (‘parton level’ refers to the b -quarks produced by the Higgs decays, before hadronisation or showering), with other variables found in appendix A. Ideally, analyses should have a high signal-to-background ratio S/B across different kinematic regimes of the Higgs boson as their boosts can vary rapidly with variations in coupling.

Figure 3a shows κ_λ values between 1 and 4 inclusive, where destructive interference is near maximal and small κ_λ variations cause dramatic changes to the m_{hh} and $p_{\text{T}}(h)$ distributions. For SM couplings, events are suppressed at low values $250 < m_{hh} \lesssim 300 \text{ GeV}$ and peaks at $m_{hh} \simeq 400 \text{ GeV}$, resulting in higher p_{T} Higgs bosons. For $\kappa_\lambda = 2$, destructive interference causes the m_{hh} shape to vanish at $m_{hh} \simeq 320 \text{ GeV}$. The signal instead occupies localised regions on either side of this minimum: near the kinematic threshold and at higher



(a) κ_λ near maximal destructive interference (b) κ_λ constructive vs destructive interference

Figure 3. Unit normalised distributions of the (upper) di-Higgs invariant mass m_{hh} and (lower) Higgs transverse momentum $p_T(h)$ at parton level. The κ_λ values highlight differential effects (a) around the point of maximal destructive interference $\kappa_\lambda \simeq 2.5$ and (b) of constructive vs destructive interference arising from the sign of κ_λ .

values where the Higgs bosons are more boosted. As κ_λ increases to $\kappa_\lambda = 3$ and 4, many events shift to low $m_{hh} \lesssim 350$ GeV, where the Higgs bosons have comparatively lower p_T .

Figure 3b shows how the sign of κ_λ impacts the m_{hh} and $p_T(h)$ distributions. For scenarios with destructive interference $\kappa_\lambda > 0$, the signals $\kappa_\lambda \geq 5$ all occupy low m_{hh} values peaking near 250 GeV. Interestingly, as κ_λ increases from 5 to 10, the m_{hh} and $p_T(h)$ spectra become slightly harder. Meanwhile, scenarios with constructive interference $\kappa_\lambda < 0$ have a greater proportion of signal events occupying comparatively higher m_{hh} and $p_T(h)$ values. These qualitatively different kinematic features can help lift degeneracies in the sign and value of λ_{hhh} that give the same total cross-sections.

Physically, these features arise due to kinematic thresholds of the interfering amplitudes. The triangle amplitude tends to dominate at lower m_{hh} , while the box diagram has a kinematic threshold at approximately twice the top mass $2m_{\text{top}}$ and impacts larger m_{hh} values. Reconstructing these differential features, especially where the cross-section gradient is small, can improve λ_{hhh} sensitivity.

Process	$p_T(j_1^{\text{gen}})$ [GeV]	$\sigma_{\text{LO}}^{\text{MG}}$ [pb]	$N_{\text{events}}^{\text{gen}} (\times 10^6)$	\mathcal{L}_{eff} [fb $^{-1}$]	k-factor
$2b2j$	[20, 200]	2.3×10^4	2	0.087	1.3
$2b2j$	[200, 500]	1.5×10^3	2	1.3	
$2b2j$	[500, 1000]	35.3	2	56.7	
$2b2j$	> 1000	0.706	1.1	1560	
$4b$	[20, 200]	63.2	2	31.7	1.6
$4b$	[200, 500]	2.82	2	710	
$4b$	[500, 1000]	0.041	2	4.9×10^4	
$4b$	> 1000	5.5×10^{-4}	2	3.6×10^6	
$t\bar{t}$	—	532	2.1	3.95	1.4
$t\bar{t} + b\bar{b}$	—	2.7	1.0	370	—
$t\bar{t}h$	—	0.44	1.0	2300	—
$b\bar{b}h$	—	0.076	1.0	1.3×10^4	—
ZZ	—	11.5	1.0	87	—
Zh	—	0.72	1.0	1400	—
Wh	—	1.4	1.0	710	—
hh signal (κ_λ, κ_t)					
$(1, 1)_{\text{SM}}$	—	0.016	0.1	6200	2.4
$(2, 1)$	—	0.0076	0.1	13000	2.3
$(3, 1)$	—	0.0082	0.1	12000	2.7
$(5, 1)$	—	0.037	0.1	2700	3.1
$(10, 1)$	—	0.26	0.1	380	3.1

Table 1. Summary of Monte Carlo samples. The columns indicate any explicit parton level generator preselection on the leading parton $p_T(j_1^{\text{gen}})$, the leading order MADGRAPH cross-section $\sigma_{\text{LO}}^{\text{MG}}$, the number of events generated $N_{\text{events}}^{\text{gen}}$, and an estimate of the effective luminosity $\mathcal{L}_{\text{eff}} = N_{\text{events}}^{\text{gen}} / \sigma_{\text{LO}}^{\text{MG}}$. The k-factor is a scaling factor applied to account for NLO corrections for the dominant backgrounds, and overall higher order corrections for the signal described in the main text. Information for a representative set of signals are displayed in the bottom five rows. For the signals, higher statistics samples of 250k events per point are generated for neural network training as discussed in the main text.

2.2 Background processes

We calculate cross-sections and produce MC events at LO for background processes using the same generator configuration as the signals, which are summarised in table 1. We refer to processes with four real b -quarks at parton level as ‘irreducible’ backgrounds. Processes with fewer than four real b -quarks can enter the analysis selection if a jet is misidentified as a b -jet, which we refer to as ‘reducible’ backgrounds. We apply k-factors from ref. [46] to account for NLO corrections for the dominant backgrounds shown in table 1.

We generate two sets of multijet processes: irreducible $b\bar{b}b\bar{b}$ ($4b$) and reducible $b\bar{b}j\bar{j}$ ($2b2j$). We define light-flavour partons j as all quarks and gluons except bottom and top

$j \in \{u, d, s, c, g\}$. The $2b2j$ cross-section is two orders of magnitude larger than $4b$, so is substantial even with powerful light jet rejection from b -tagging algorithms. We do not simulate $4j$ to preserve computational resources and ref. [46] showed that this process is subdominant compared to $4b$ and $2b2j$. To improve statistics in high p_T tails, we generate events for ranges of leading parton $p_T(j_1^{\text{gen}})$ in MADGRAPH, sliced with lower bin edges at 20, 200, 500, 1000 GeV, where leading parton refers to b -quarks or light partons j .

We generate top quark pairs $t\bar{t}$ and the irreducible $t\bar{t}+b\bar{b}$ process at the matrix element level, which together are expected to comprise around 10% of the background rates. For single Higgs processes in associated production, we consider $Wh, Zh, t\bar{t}h, b\bar{b}h$. These processes typically have cross-sections that are an order of magnitude greater than di-Higgs, and the presence of a Higgs boson increases the probability of events passing analysis selections than other electroweak processes. Finally, we consider the diboson ZZ process, which has a modest cross-section and its $4b$ final state constitutes an irreducible background. We generate at least one million events for each of these processes.

2.3 Detector emulation

For all signal and background samples, the decay, parton shower, hadronisation, and underlying event are modelled by PYTHIA 8.230 [131]. To save computational resources for our simplified study, we do not emulate pileup. We expect that detector upgrades [132–134] and improvements in pileup mitigation techniques [135–139] will reduce the impact of pileup. To emulate reconstruction effects, we use DELPHES 3.4.1 [140] and assume the default ATLAS configuration card unless stated otherwise. We define three sets of reconstructed jets using the anti- k_t clustering algorithm [141, 142] with radius parameter R :

- *Small jets* (j_S) are defined with $R = 0.4$ and cluster only calorimeter towers. We impose $p_T > 40$ GeV on all these jets to emulate the detector trigger requirement. We require these to be within a tracking acceptance of $|\eta| < 2.5$ to allow b -tagging, where η is the pseudorapidity.
- *Large jets* (j_L) are defined with $R = 1.0$, also clustering only calorimeter towers. Based on the expected kinematics of boosted Higgs bosons $p_T \gtrsim 2m_h$, we require these to satisfy $p_T > 250$ GeV and be central $|\eta| < 2.0$.
- *Track jets* (j_T) are defined by clustering only tracking information using $R = 0.2$. We impose kinematic requirements of $p_T > 20$ GeV and $|\eta| < 2.5$. We associate these track jets to large jets if their distance satisfies $\Delta R(j_T, j_L) < 1.0$, where $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$.

For simplicity, we tag the flavour of both small jets and track jets using the same default implementation in DELPHES, which parameterises tagging efficiencies based on the truth quark flavour of the jet. The p_T -dependent efficiencies of bottom, charm, and light jet (u, d, s, g) mistag rates are based on the ‘70% working point’ of the ATLAS multivariate MV2c20 algorithm [71]. The b -tagging efficiency peaks at 74% for $p_T \simeq 150$ GeV, falling to 50% at $p_T \simeq 500$ GeV. The charm mistag rate is around 10%, peaking at 14% for

$p_T \simeq 100$ GeV before falling to 7% at $p_T \simeq 500$ GeV. The light mistag rate remains on the order of 1% throughout the p_T regimes of interest. We do not increase the b -tagging range to $|\eta| < 4.0$ expected from the detector upgrades [143] because we expect the region $|\eta| > 2.5$ to be dominated by the gluon splitting background $g \rightarrow b\bar{b}$ whereas b -quarks from Higgs are more central. However, we follow ref. [57] to emulate expected b -tagging improvements of 8% per jet in $|\eta| < 2.5$ by scaling the $4b$ background and hh signals by a factor of 1.36, and the $2b2j$ and $t\bar{t}$ backgrounds by 1.17. To improve MC statistics after b -tagging, we multiply events by these efficiencies as a weight rather than discarding non- b -tagged events. For processes with a low number of real b -quarks such as $2b2j$, this can improve the number of raw MC events by up to three orders of magnitude when requiring four b -tagged jets.

Electrons and muons must be within the tracking acceptance $|\eta| < 2.5$, satisfy $p_T > 10$ GeV, together with default DELPHES efficiencies, smearing and isolation. We modify the default calculation of missing transverse momentum ($\mathbf{p}_T^{\text{miss}}$) to ensure muons are included.

It is relevant to study how experimental effects such as jet clustering, missing energy due to neutrinos and finite detector resolution impact the discriminating variable m_{hh} . As m_{hh} is sensitive to κ_λ , this can motivate improvements to its reconstruction for future work. Figure 4 shows the impact of these effects for the $\kappa_\lambda = 2.5$ signal with near-maximal destructive interference to accentuate the shape differences. First, four small jets (with the p_T and η requirements defined above applied) are identified by association to the b -quarks from the Higgs bosons found in the truth record. For these jets, we then compare their parton level m_{hh} distribution (shaded grey) to three classes of jets at different levels of reconstruction: 1) truth-level jets clustered from *all* final state truth particles (yellow line), 2) truth-level jets clustered from all final state truth particles *except neutrinos* (blue line), and 3) reconstructed-level jets (red line). The shape of m_{hh} for all three jet definitions is mildly distorted to lower values than the parton level, with the reconstructed jets being generally lower than their truth counterparts. While this shows the degradation of m_{hh} resolution, the qualitative features of the non-trivial m_{hh} shape are preserved.

Figure 4a considers a $p_T > 20$ GeV threshold on the reconstructed jets, while figure 4b raises this threshold to the nominal $p_T > 40$ GeV considered in this work. This illustrates that the signal rate at low m_{hh} is substantially driven by trigger thresholds. Future work may also consider jets with variable radius, which may improve signal-to-background ratios [144]. This comparison highlights the following important experimental considerations:

- Maintaining sufficiently low trigger thresholds for the HL-LHC upgrades [132] is of key importance for both discovery of the di-Higgs process as well as constraining λ_{hhh} . In particular, figure 3a shows the $1 \leq \kappa_\lambda \leq 4$ scenarios exhibit large qualitative changes in shape at low $m_{hh} \lesssim 400$ GeV. A rise in trigger thresholds can reduce sensitivity to this region. Figure 3b shows the majority of $\kappa_\lambda \geq 5$ signals reside at $m_{hh} \lesssim 350$ GeV, but this is also compensated by the faster changes in total cross-section.
- Corrections should be applied to compensate for energy loss in b -jets, which is not accounted for by standard jet calibration. These techniques have been recently deployed by the LHC collaborations [145] but are beyond the scope of this work. Applying these corrections at the trigger level remains an open problem.

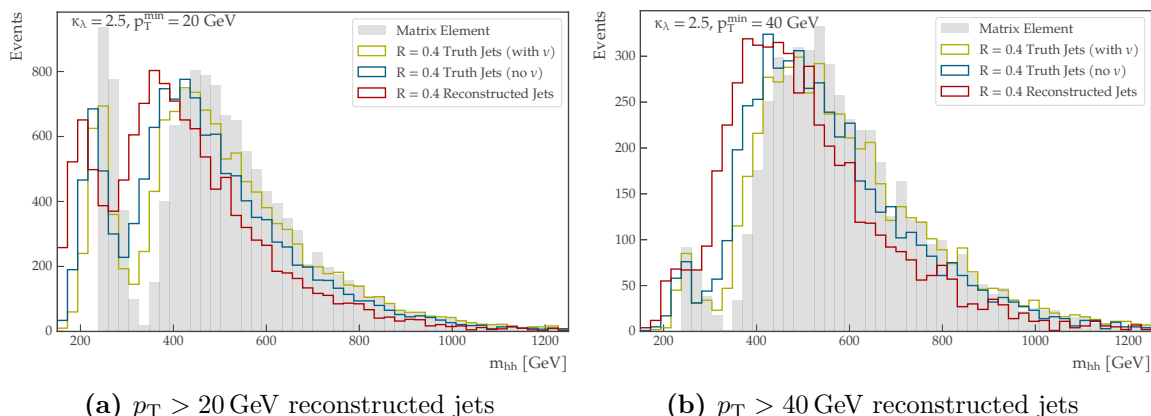


Figure 4. Distributions of the di-Higgs invariant mass m_{hh} for $\kappa_\lambda = 2.5$ where destructive interference is near-maximal. These are shown at different levels of reconstruction: the parton level predictions (grey shaded), truth jets including neutrinos (yellow line), truth jets without including neutrinos (blue line), and reconstructed jets after the detector emulation described in subsection 2.3 (red line). Small $R = 0.4$ radius jets are used that satisfy $|\eta| < 2.5$. The p_T thresholds in the sub-figure caption are applied to the reconstructed jets.

3 Analysis strategies

This section presents our analysis strategies to constrain the trilinear self-coupling λ_{hhh} in the $hh \rightarrow 4b$ channel. We first outline the figures of merit that motivate our analysis design. The goal of measurement is to maximise discrimination between different (λ_{hhh}, y_t) values, which requires solving two conceptually distinct classification problems:

- *signal characterisation* i.e. the measurement power to discriminate λ_{hhh}^i vs $\lambda_{hhh}^{j \neq i}$. This is most simply quantified by the difference (squared) of the signal rates $(S_i - S_{j \neq i})^2$.
- *background suppression* i.e. signal S vs background B discrimination. Intuitively, if the changes $(S_i - S_{j \neq i})^2$ between two couplings are comparable or smaller than the background uncertainties, this reduces the ability to discriminate λ_{hhh}^i vs $\lambda_{hhh}^{j \neq i}$.

We quantify how well our analyses simultaneously achieves these goals in a statistically meaningful way using the chi-square

$$\chi_{ij}^2 = \frac{(S_i - S_j)^2}{\varsigma_{S_i}^2 + \varsigma_B^2}. \quad (3.1)$$

Here, S denotes the signal yields after analysis selections for the nominal (λ_{hhh}^i, y_t^i) and alternative (λ_{hhh}^j, y_t^j) coupling hypotheses. The ς_S (ς_B) are the combined absolute uncertainties on the signal (background), where $\varsigma_B \gg \varsigma_{S_i}$ is typical in the $4b$ channel.

To benchmark SM measurement precision, we can fix $S_j = S_{\text{SM}}$, which assumes the observed data will correspond to that of the SM. Nonetheless, it is important to consider dedicated optimisation assuming $S_i = S_{\text{BSM}}$ should nature prefer BSM couplings. Higher χ^2 values indicate better discrimination power between two coupling hypotheses, which is achieved by maximising the numerator $(S_i - S_j)^2$ while minimising the uncertainties $\varsigma_{S,B}$.

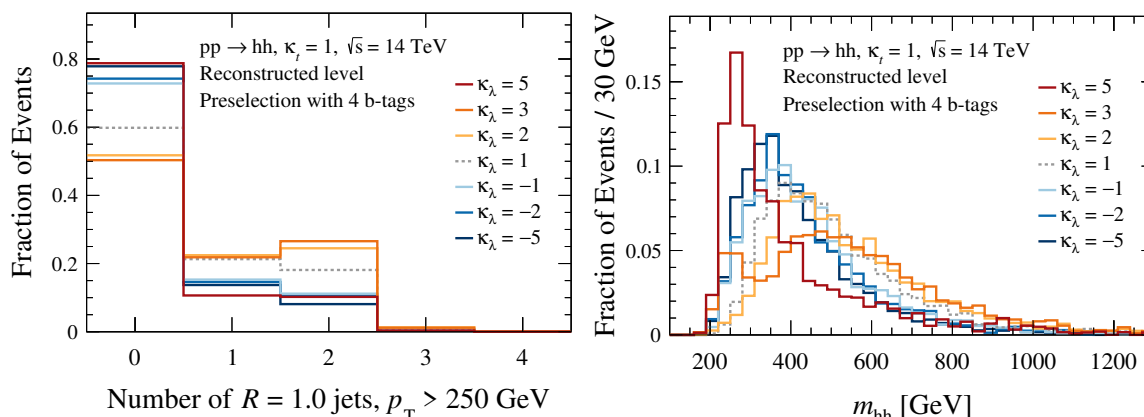


Figure 5. Unit normalised distributions of $pp \rightarrow hh$ signals with different κ_λ (solid lines) compared with the SM (grey dashed) at reconstructed level. The preselection with four b -tagged jets is applied. Displayed are (left) the number of reconstructed large jets and (right) the m_{hh} variable. Negative (positive) values of κ_λ are coloured by shades of blue (orange). Large jets have a radius parameter of $R = 1.0$ and $p_T > 250$ GeV.

As background rates are large, the systematic component of the uncertainty dominates the denominator. One way to reduce the impact of background systematics is to suppress B . We design two classes of analysis strategies to fulfil these objectives:

- The *baseline analysis* (subsection 3.1) uses conventional rectangular cuts on variables reconstructing Higgs bosons, inspired by recent ATLAS and CMS strategies. This serves as a baseline to benchmark the performance of neural network optimisation.
- The *neural network analysis* (subsection 3.2) demonstrates the use of a multivariate strategy to optimise sensitivity beyond the baseline analysis. This consists of an additional selection requirement based on the output of an artificial neural network.

3.1 Baseline analysis

The *baseline analysis* is loosely inspired by a recent ATLAS analysis [41]. Our event selection is summarised in table 2. In all categories, we require at least four b -tagged small or track jets to suppress multijet backgrounds. We also require the pseudorapidity difference of the two reconstructed Higgs candidates to be small $|\Delta\eta(h_1, h_2)| < 1.5$ because high mass objects occupy more central regions than multijet backgrounds dominated by gluon-gluon scattering. To suppress $W \rightarrow \ell\nu$ decays from top quarks, we veto electrons or muons and require $E_T^{\text{miss}} < 150$ GeV in all categories.

To probe different regimes of Higgs boson kinematics, we define three categories based on the exclusive number of large jets in each event, which we refer to as *resolved*, *intermediate*, and *boosted*. Figure 5 displays the multiplicity of resolved, intermediate and boosted events for different κ_λ , as well as reconstructed m_{hh} distributions inclusive of number of large jets. The multiplicity distribution shows that this categorisation has λ_{hhh} discrimination. Interestingly, $\kappa_\lambda = 2$ and $\kappa_\lambda = 3$ have more events falling in the intermediate and boosted categories $N(j_L) \geq 1$ than the SM $\kappa_\lambda = 1$, consistent with the higher m_{hh} and

$p_T(h)$ tails (figure 3a). This orthogonal categorisation in $N(j_L)$ enables straightforward statistical combination to enhance sensitivity. For the three categories, the final signal region (SR) is defined by mass window requirements on the reconstructed $h \rightarrow bb$ systems of $m(h_{1,2}) \in [90, 140]$ GeV, based on the mass resolution and background rejection.

Resolved. The resolved category requires exactly zero reconstructed large jets. This targets Higgs bosons with low transverse momenta, $p_T \lesssim 2m_h$, which we reconstruct as four distinct small jets. To identify pairs of jets consistent with a $h \rightarrow bb$ decay, we consider the leading four small jets and construct Higgs boson candidates from small jet pairs that minimise the mass difference between the candidates $\Delta m(h_1, h_2)$. The (sub)leading Higgs candidate is defined as the system of a pair of small jets with the (lower) higher p_T . As the triangle amplitude dominates at low m_{hh} and therefore low Higgs p_T , this category is particularly important for λ_{hhh} sensitivity. We implement a selection defined in ref. [41], where the angular distance $\Delta R_{jj}^{h_{1,2}}$ between the jet pair of the Higgs candidates satisfy

$$m_{4j} < 1250 \text{ GeV} : \begin{cases} -0.5 + \frac{360 \text{ GeV}}{m_{4j}} < \Delta R_{jj}^{h_1} < \frac{653 \text{ GeV}}{m_{4j}} + 0.475, \\ \frac{235 \text{ GeV}}{m_{4j}} < \Delta R_{jj}^{h_2} < \frac{875 \text{ GeV}}{m_{4j}} + 0.35, \end{cases} \quad (3.2)$$

$$m_{4j} \geq 1250 \text{ GeV} : \Delta R_{jj}^{h_{1,2}} < 1. \quad (3.3)$$

This adjusts the angular distance between the jets of each Higgs candidate according to the boost of the system characterised by the invariant mass of the 4-jet system m_{4j} .

Intermediate. The intermediate category requires exactly one large jet in the event. This targets regimes where exactly one Higgs boson is sufficiently boosted ($p_T \gtrsim 2m_h$) that the b -jets in the $h \rightarrow bb$ system become merged so are more efficiently reconstructed as one large jet. The two b -quarks are reconstructed as two track jets j_T , which we require to be associated to this large jet by $\Delta R(j_T, j_L) < 1.0$. The remaining small jets j_S , separated from the large jet by $\Delta R(j_S, j_L) > 1.2$, are paired to form the subleading Higgs candidate, where the j_S pair is chosen to minimise the mass difference of the two Higgs candidates.

Boosted. The boosted category requires exactly two large jets targeting $hh \rightarrow 4b$ events where both Higgs bosons have high Lorentz boosts. These events typically reside in the tails of the m_{hh} distribution ($m_{hh} \gtrsim 500$ GeV), which can be important for probing BSM λ_{hhh} couplings that enhance m_{hh} and $p_T(h)$ at high values. Importantly, the high expected signal rate due to the large branching ratio $\mathcal{B}(h \rightarrow bb) \simeq 58\%$ implies greater statistical power in these tails compared with lower rate channels such as $bb\tau\tau$ or $bb\gamma\gamma$. We require each large jet to have two b -tagged track jets associated to this large jet.

Figure 6 shows the signal acceptance times efficiency $A \times \varepsilon$ for the *baseline analysis* signal region in the three categories. Due to the different Higgs boson kinematics when varying κ_λ , the $A \times \varepsilon$ strongly depends on κ_λ . We find the highest $A \times \varepsilon = 0.55\%$ at $\kappa_\lambda = 1.5$ in the resolved category while the largest $A \times \varepsilon = 0.19\%$ (0.21%) for the intermediate (boosted) category peaks at $\kappa_\lambda = 2$, where destructive interference is near maximal and the Higgs bosons have the greatest boost (figure 3a). The $A \times \varepsilon$ values in the intermediate and boosted categories decrease precipitously outside $1 \lesssim \kappa_\lambda \lesssim 3$ and become an order

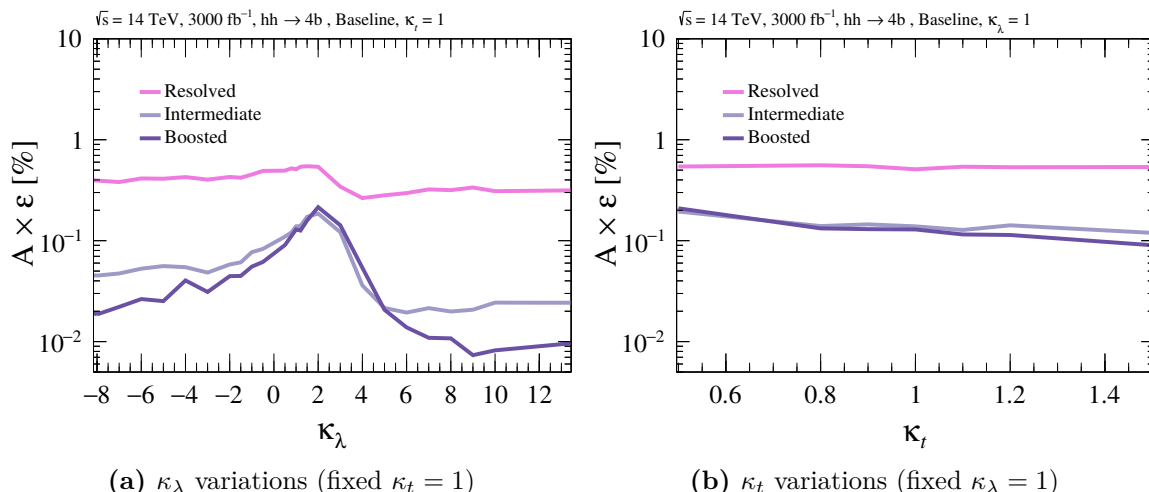


Figure 6. Signal acceptance times efficiency $A \times \varepsilon$ in percent for the *baseline analysis* for variations in (a) κ_λ and (b) κ_t . This is shown for the resolved (pink), intermediate (lilac) and boosted (purple) categories. The $A \times \varepsilon$ is equivalent to the number of signal events S divided by initial number of events $\sigma \times \mathcal{L}$.

of magnitude lower than those of the resolved category. In particular, we find the lowest $A \times \varepsilon = 0.26\%$, 0.019% , 0.007% at $\kappa_\lambda = 4, 6, 9$ for the resolved, intermediate, and boosted categories respectively. This suppression is dominantly due to the lower $h \rightarrow bb$ boosts causing events to pass the jet (trigger) p_T requirements with lower efficiency. Figure 6b shows that the acceptance decreases slowly for increasing κ_t due to the m_{hh} distribution being shifted toward lower values, as displayed in appendix A.

Figure 7 shows the m_{hh} and leading Higgs p_T distributions for signal and background in the signal region of the three categories. Notably, the resolved category has the greatest shape discrimination between κ_λ hypotheses, where $\kappa_\lambda = 5$ has a visibly higher proportion of events at lower m_{hh} values than $\kappa_\lambda = 1$. To exploit this, our final selection divides events into non-overlapping m_{hh} bins whose lower edges are defined in table 2. The resolved category uses six bins chosen for simplicity, while fewer bins are used for the intermediate and boosted categories due to lower expected event rates. When performing the statistical analysis in section 4, these orthogonal m_{hh} bins are combined to enhance sensitivity. Future work could consider dedicated optimisation of this binning scheme. Appendix A shows further distributions involving the subleading Higgs p_T and di-Higgs system $p_T(hh)$.

3.2 Neural network analysis

Signal characterisation (λ_{hhh}^i vs $\lambda_{hhh}^{j \neq i}$) and background suppression (S vs B) are demanding classification problems seeing promising adoption of deep neural network (DNN) solutions in particle physics [77–87]. We implement our *neural network analysis* using a single network architecture with the KERAS library [146]. This comprises a feed-forward network [74–76] with a depth of two internal (‘hidden’) layers each with 200 nodes, similar to ref. [147], densely connected to each other and to the input and output nodes. The internal

Observable	Preselection		
Large jet j_L	$R = 1.0, p_T > 250 \text{ GeV}, \eta < 2.0$		
Small jet j_S	$R = 0.4, p_T > 40 \text{ GeV}, \eta < 2.5$		
Track jet j_T	$R = 0.2, p_T > 20 \text{ GeV}, \eta < 2.5$		
$j_T \in j_L$	$\Delta R(j_T, j_L) < 1.0$		
	Resolved	Intermediate	Boosted
$N(j_L)$	$= 0$	$= 1$	$= 2$
$N(j_S)$	≥ 4	≥ 2	≥ 0
h_1^{cand}	$j_S^{(i)}$ pair	j_L	$j_L^{(1)}$
h_2^{cand}	$j_S^{(i)}$ pair	$j_S^{(i)}$ pair, $\Delta R(j_S^{(i)}, j_L) > 1.2$	$j_L^{(2)}$
ΔR_{jj}	See eqs. (3.2), (3.3)	—	—
Signal region			
$j_T \in h_1^{\text{cand}}$	—	≥ 2	≥ 2
$j_T \in h_2^{\text{cand}}$	—	—	≥ 2
b -tagging	Two b -tags for each h_i^{cand}		
$ \Delta\eta(h_1, h_2) $	< 1.5		
E_T^{miss}	$< 150 \text{ GeV}$		
$p_T^\ell, \eta_\ell $	$> 10 \text{ GeV}, < 2.5$		
N_ℓ	$= 0$		
$p_{\text{signal}}^{\text{DNN}}$	> 0.75 (<i>neural network analysis</i> only)		
	Resolved	Intermediate	Boosted
$m(h_1) \text{ [GeV]}$	[90, 140]	[90, 140]	[90, 140]
$m(h_2) \text{ [GeV]}$	[90, 140]	[90, 140]	[90, 140]
Lower bin edges for m_{hh} binning [GeV]			
Resolved	[200, 250, 300, 350, 400, 500]		
Intermediate	[200, 500, 600]		
Boosted	[500, 800]		

Table 2. Overview of event selection for the *baseline analysis* in the resolved, intermediate and boosted categories. The requirements above the ‘Signal region’ heading are the same as the preselection used for the *neural network analysis* training. The requirements below the ‘Signal region’ heading are the signal region requirements. The lower bin edges for the m_{hh} binning scheme is shown in the bottom three rows. See the main text for details.

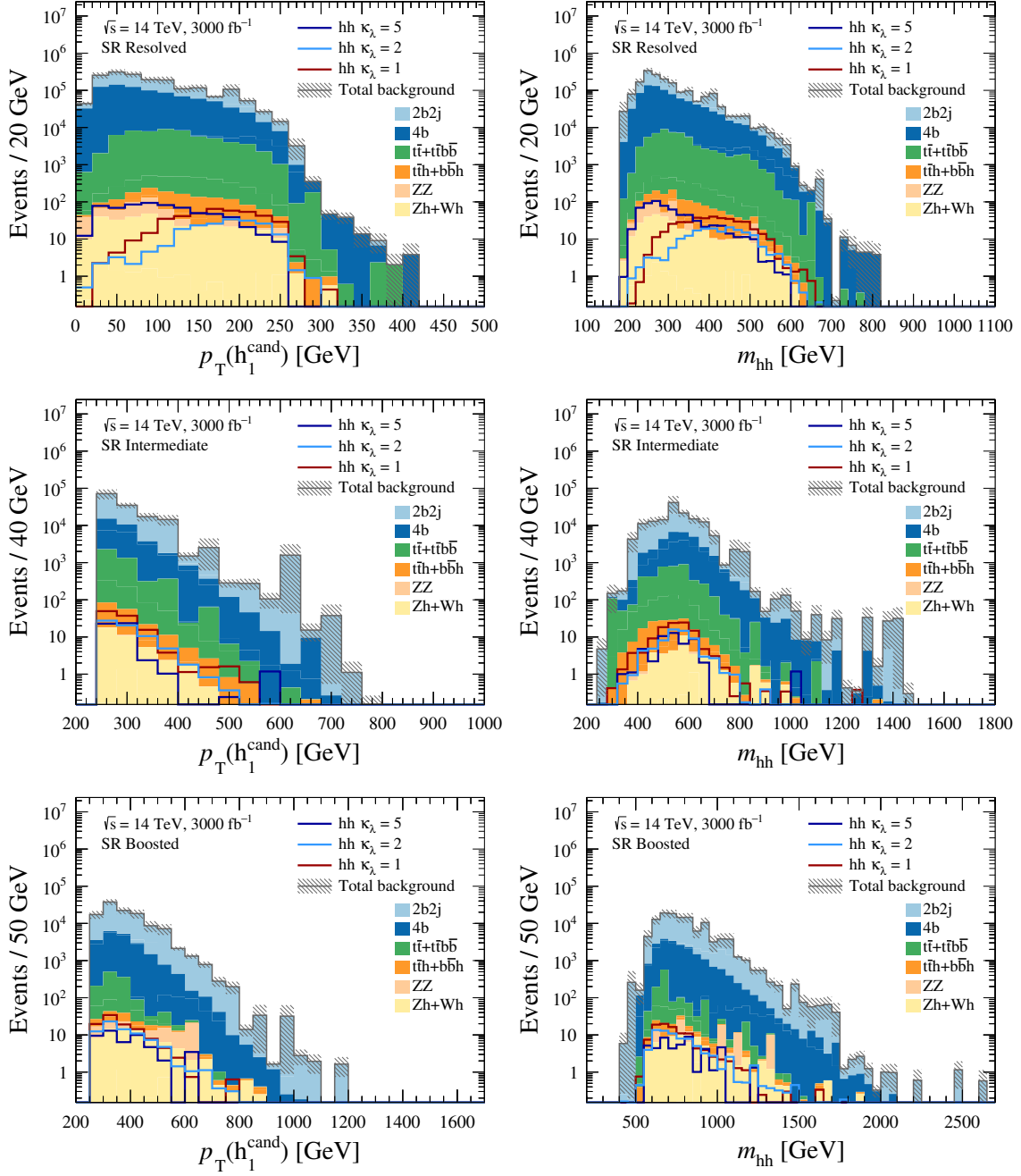


Figure 7. The *baseline analysis* leading Higgs candidate transverse momentum $p_T(h_1)$ (left) and invariant mass of the pair of Higgs boson candidates m_{hh} (right), for benchmark signals (lines) and background (filled). These are displayed for resolved (upper), intermediate (middle) and boosted (lower) categories after all kinematic selections, including $m(h_i)$ mass window cuts, are applied.

Reconstructed objects	Variables used for training
Higgs candidates $h_{1,2}^{\text{cand}}$	(p_T, η, ϕ, m)
Subjets $\in h_{1,2}^{\text{cand}}$	$\Delta R(j_1, j_2)$
Missing transverse momentum	$E_T^{\text{miss}}, \phi(\mathbf{p}_T^{\text{miss}})$
Leptons	N_e, N_μ
b -tagging	Boolean for $j_i \in h_{1,2}^{\text{cand}}$
Di-Higgs system	p_T^{hh}, m_{hh}

Table 3. Input variables used to train the neural network.

nodes use the rectified linear unit (ReLU) defined as $\max(0, x)$ for the activation function, whose advantages over the traditionally used sigmoid function are discussed in ref. [148].

As input, the network uses a comprehensive set of 20 variables summarised in table 3. This comprises the four-momenta of the two Higgs candidates, the ΔR distance between the two subjets associated to each Higgs candidate, the b -tagging state of these subjets, the missing transverse momentum with magnitude E_T^{miss} and azimuthal angle ϕ , the number of reconstructed electrons and muons, and the mass and transverse momentum of the di-Higgs system. These variables are chosen for their signal vs background discrimination power. A separate neural network is trained for each of the resolved, intermediate and boosted categories using the preselection defined in table 2. For the training, the input samples are normalised such that equal weight is given to signal, multijet ($2b2j$ and $4b$) background, and $t\bar{t}$ background.

For outputs, we construct three nodes corresponding to signal, multijet ($2b2j$ and $4b$) background, and $t\bar{t}$ background. This is referred to as a ‘multi-class classification network’. The model assigns a score p_i corresponding to how likely an event corresponds to one of the three processes. The output nodes use a normalised exponential activation function known as ‘softmax’, which is a standard choice for multi-class configurations. This constrains each output score² to $p_i \in [0, 1]$ and their sum to unity $\sum_i p_i = 1$. We constructed a three-output classifier to explore its utility in classifying background processes for designing control regions, but section 4 will only use the binary signal classifier for simplicity. In typical experimental implementations, multijet processes are estimated using data-driven methods while $t\bar{t}$ employs MC, which have different systematics.

Half of the MC events for the multijet and $t\bar{t}$ background samples, and all of the high statistics $hh \rightarrow 4b$ signal samples are set aside for training (20% of the training events are used for training validation and are not used for actual training). The categorical cross-entropy is used as the loss function, which quantifies the accuracy of the model predictions at each training step, and is minimised using the ADAMAX algorithm [149]. The step size used in this minimisation is controlled by a *learning rate* hyperparameter. The cross-entropy H between the network prediction and the true classes in N events from the training set is defined by $H(p^{\text{label}}, p^{\text{model}}) = -\frac{1}{N} \sum_{i=1}^N p_i^{\text{label}} \log p_i^{\text{model}}$. Here, p^{label} is

²Despite the notation and its properties, p_i is not a true posterior probability.

the vector containing the class of each event (1 for the true class and 0 for the other two) and p^{model} is a vector containing the score assigned to each class for each event. The initial learning rate in the ADAMAX optimiser is set to 5×10^{-3} for the resolved and intermediate neural networks, and 5×10^{-5} for the boosted ones. The training set is divided into batches during training. Once all batches are finished processing, one training epoch is complete and the next one starts by processing the first batch again. The number of events in each batch is a tunable hyperparameter, which we set to 100 in this study. To mitigate overfitting unphysical features such as statistical fluctuations, we apply dropout [150] at a rate of 30% to both internal layers. This means 30% of internal nodes are randomly masked during each training iteration. We find 20 training epochs gives close to optimal performance. The learning rate, batch size and dropout rate are optimised using a random search method.

Figure 8 shows the signal score $p_{\text{signal}}^{\text{DNN}}$ distributions for the DNN trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$ for background and benchmark signals in the three categories. The signal vs background discrimination is improved across the categories, suggesting that our neural networks capture kinematic information beyond the cuts of the *baseline analysis*. However, this depends on the value of κ_λ . For example, the upper-left plot shows that the DNN trained on $\kappa_\lambda = 1$ adds substantial discrimination power for a $\kappa_\lambda = 1$ signal, but not for a $\kappa_\lambda = 5$ signal. This will be further discussed in section 4. Our *neural network analysis* imposes a universal requirement of $p_{\text{signal}}^{\text{DNN}} > 0.75$ for simplicity. Future work could consider optimising requirements on $p_{\text{signal}}^{\text{DNN}}$ that are different for each category. Additionally, one could extend the analysis by fitting the $p_{\text{signal}}^{\text{DNN}}$ variable instead of using only one bin.

Figure 9 shows the m_{hh} distributions after the $p_{\text{signal}}^{\text{DNN}} > 0.75$ requirement is imposed for the DNN trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$. We note that the shape of the $\kappa_\lambda = 5$ signal (as well as that of the background) in the resolved category is particularly different between the two trainings. Since the low- m_{hh} regime offers the most discrimination between different κ_λ values, we use $\kappa_\lambda = 5$ as the nominal signal training sample to ensure the DNN gives more weight to these events with respect to a SM optimisation. As we will see in section 4, this indeed performs better than training on $\kappa_\lambda = 1$ when setting κ_λ coupling limits. Note that the same is not necessarily true when setting cross-section limits assuming SM kinematics.

3.2.1 What is the DNN learning?

As machine learning techniques become increasingly widespread, it is important to understand how our neural network exploits the given physics information [82]. Specifically, we evaluate the ranked *feature importance* to quantify how much each input variable (model *feature*) changes the signal score in both signal and background events. Interpretability of deep neural networks has seen rapid development in the computational sciences. We adopt a recently developed framework for interpreting predictions called SHapley Additive exPlanations (SHAP) [88]. This framework combines several feature importance tests available for machine learning models in the literature into a single value as detailed in ref. [88]. These values are consistent for all types of inputs to the neural network, and can be compared with one another. A heuristic description of SHAP values can be found in appendix B.1 that provides intuition for this approach.

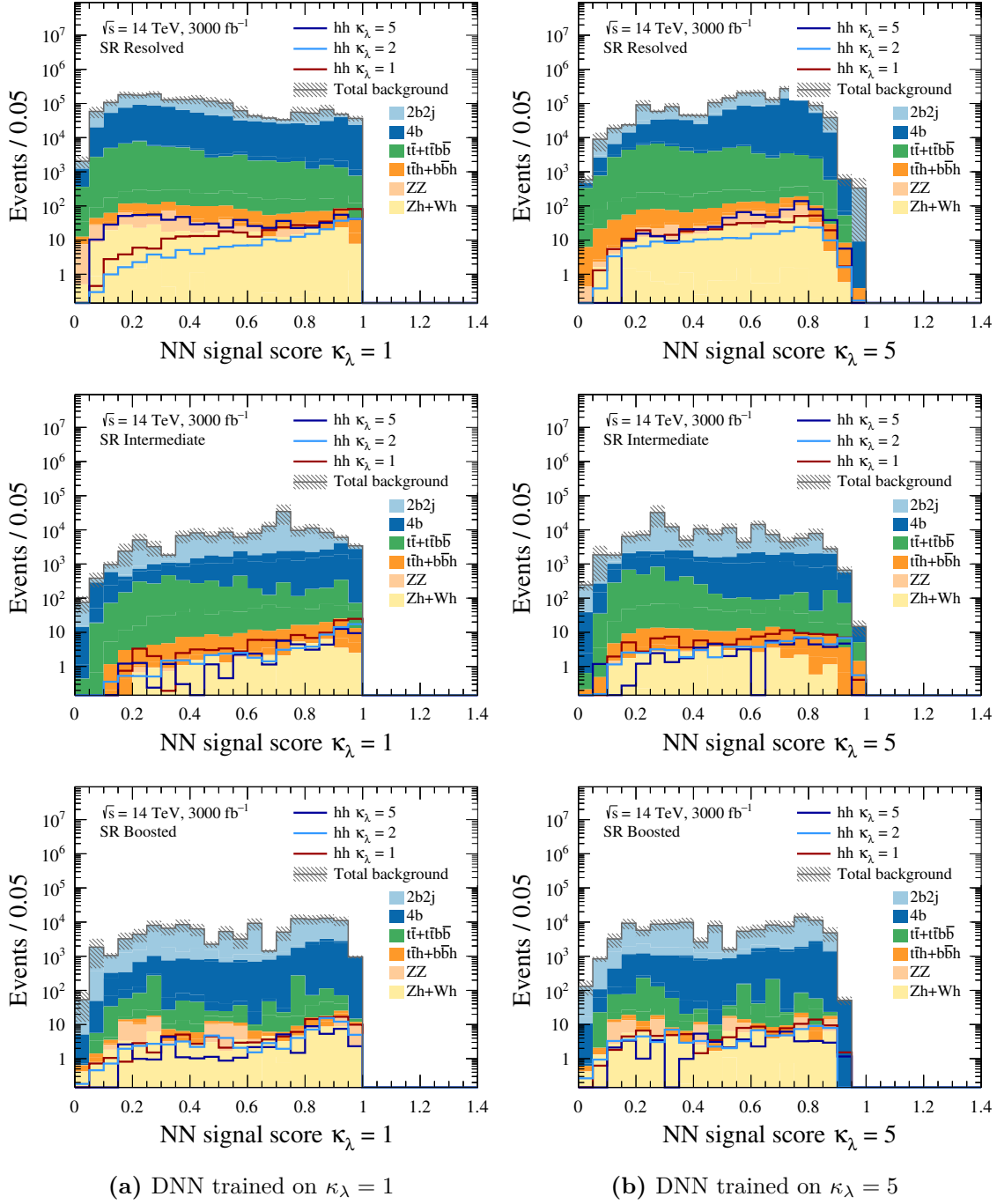


Figure 8. Neural network score distributions $p_{\text{signal}}^{\text{DNN}}$ of benchmark signals (solid lines) and background processes (filled stacked) displayed in the legend. All event selection criteria of the *neural network analysis* except the $p_{\text{signal}}^{\text{DNN}} > 0.75$ requirement are imposed. The DNN is trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signals. These are displayed for (upper) resolved, (middle) intermediate and (lower) boosted categories. The plots are normalised to $\mathcal{L} = 3000 \text{ fb}^{-1}$.

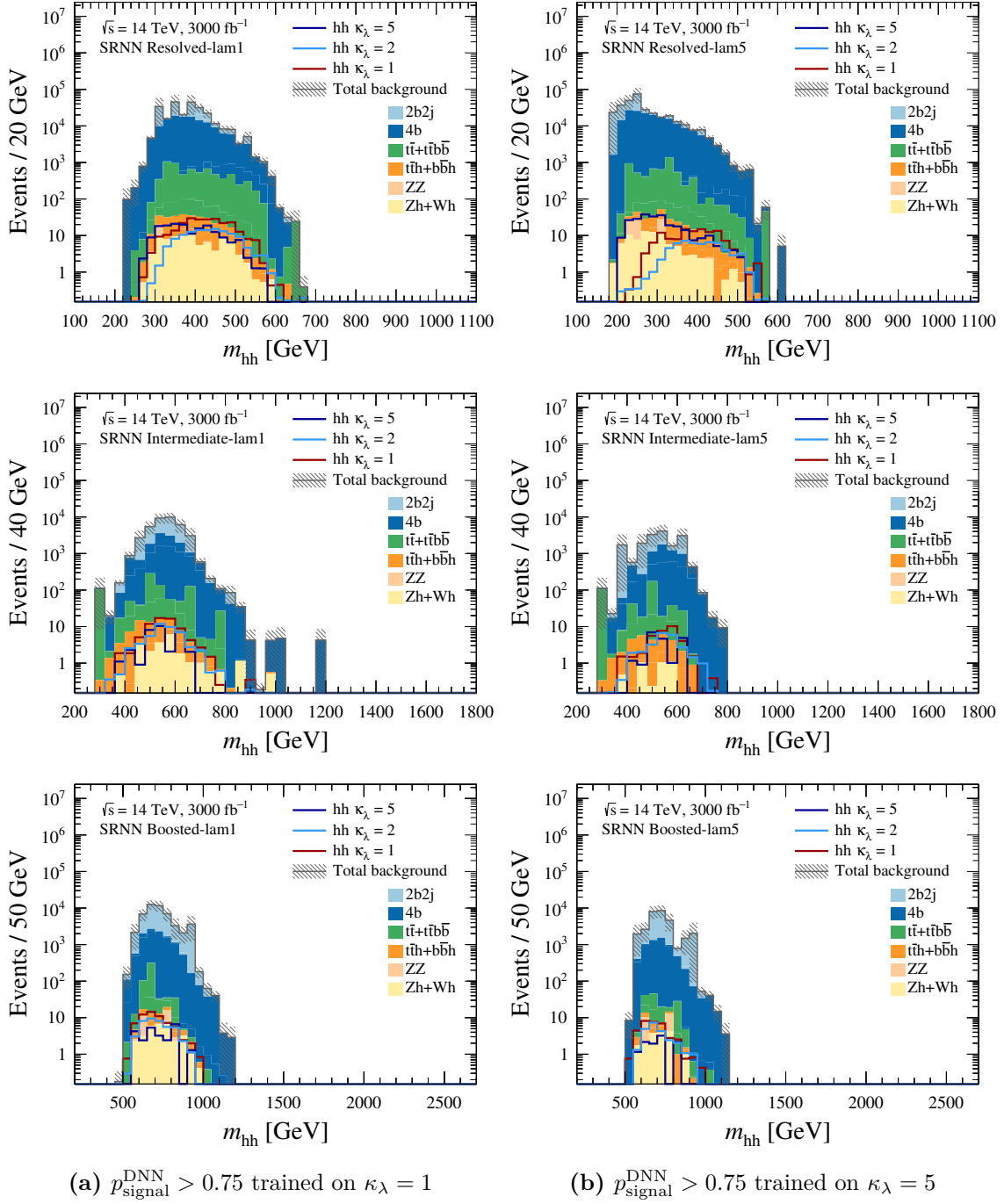


Figure 9. The m_{hh} distributions of benchmark signals (solid lines) and background processes (filled stacked) displayed in the legend. All event selection criteria of the *neural network analysis* are imposed. The DNN is trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signals. These are displayed for (upper) resolved, (middle) intermediate and (lower) boosted categories. The plots are normalised to $\mathcal{L} = 3000 \text{ fb}^{-1}$.

The SHAP framework requires that the trained model is applied to a specific set of events and only evaluates the feature importance for these events. We construct a new subset of events with half signal and half background to reflect the importance of distinguishing them from each other rather than discrimination between the background components had the S/B mirrored the preselection rates. For this evaluation, we construct the background sample to be composed of 80% $2b2j$, 15% $4b$ and 5% $t\bar{t}$ events to approximate the background composition of the three categories.

Figure 10 shows the 20 input variables of the neural networks, ranked by their impact on the final signal score for the $\kappa_\lambda = 5$ training. This impact is measured by the magnitude of their SHAP values averaged over the whole dataset given to the framework, which is plotted on the x -axis. Each point plotted per row corresponds to one event fed to the framework, and its location along the x -axis represents what impact that variable has on the signal score of the event. The relative magnitude represents how much the value of that variable changes the signal score compared to all other variables in all events. Points with a larger positive SHAP value increase the signal score more while the inverse is true for negative SHAP values. The absolute scale of the SHAP value is arbitrary in these plots. The colour scale indicates the value of the feature on the specific event e.g. a blue dot on the $b\text{-tag}(h_1^{\text{cand}}, j_1)$ indicates the leading subjet associated to the leading Higgs candidate is not b -tagged. Meanwhile, a pink dot in the m_{hh} row indicates that event has high di-Higgs invariant mass for the plotted SHAP value.

The b -tagging state of (sub-)jets are among the most important features in the resolved and intermediate categories. This is less important in the boosted category possibly due to lower b -tagging efficiencies at high p_T , which could be improved by future work in novel b -tagging techniques [151]. Angular and mass variables are stronger discriminants against multijet processes in this regime. The opening angles between these (sub-)jets and the invariant mass of the di-Higgs system carry a large amount of information in all three categories. Variables sensitive to semi-leptonic top decays, such as the number of leptons or missing transverse momentum, are effective at rejecting background (large negative SHAP value for high feature value) but less so at identifying signal (small positive SHAP values for any feature value).

Appendix B presents supplementary material characterising the performance of our neural networks. This includes receiver operator characteristic (ROC) curves, signal acceptance times efficiency, ranked feature importance for the DNN trained on $\kappa_\lambda = 1$, together with correlations between the neural network scores and reconstructed (di-)Higgs mass variables $m(h_1)$ and m_{hh} .

4 Higgs self-coupling constraints

This section presents the results and discussion for the *baseline* and *neural network* analyses. We compare the performance of the resolved, intermediate, boosted categories, and their combined constraints. We set the luminosity to the target HL-LHC dataset of $\mathcal{L} = 3000 \text{ fb}^{-1}$. Subsection 4.1 presents the signal and background rates in the signal regions of our analyses. Subsection 4.2 then performs the χ^2 statistical analysis to determine

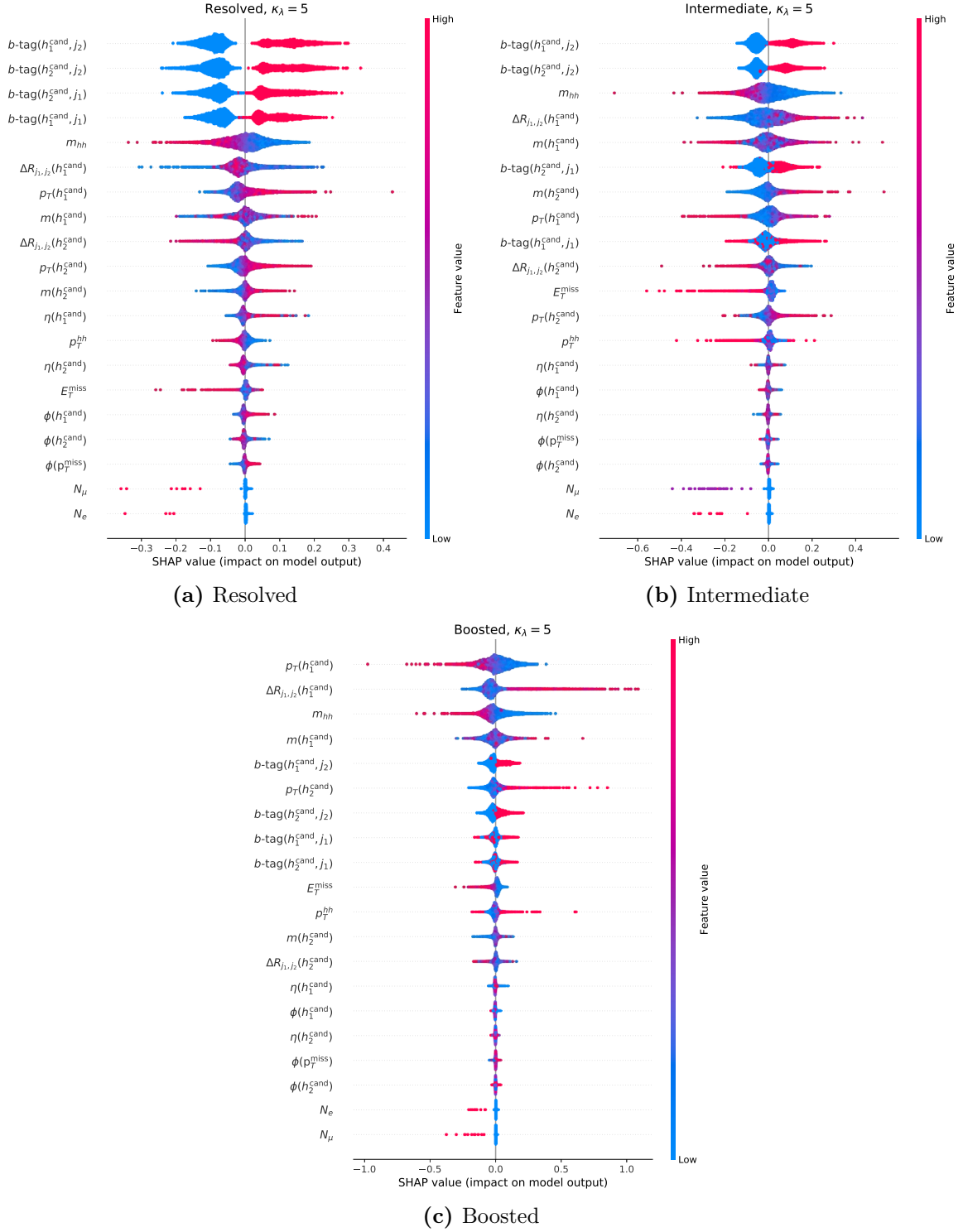


Figure 10. SHAP value plots representing ranked variable (feature) importance for the DNN trained on $\kappa_\lambda = 5$ signals in the (a) resolved, (b) intermediate and (c) boosted categories. The input variables are ranked in descending order by their average absolute SHAP value plotted on the x -axis. The colour scale indicates the value of the variable on the specific event for which the SHAP value is plotted. Analogous plots for the DNN trained on $\kappa_\lambda = 1$ are found in appendix B.

Category	Resolved			Intermediate			Boosted		
Analysis	Baseline	DNN	DNN	Baseline	DNN	DNN	Baseline	DNN	DNN
Trained on	—	$\kappa_\lambda = 5$	$\kappa_\lambda = 1$	—	$\kappa_\lambda = 5$	$\kappa_\lambda = 1$	—	$\kappa_\lambda = 5$	$\kappa_\lambda = 1$
$2b2j$	889000	134000	116000	112000	10400	24600	95300	24100	38400
$4b$	810000	183000	137000	28500	6170	13300	19700	6000	10900
$t\bar{t}$	60400	5430	6100	3650	339	639	915	76.8	385
$t\bar{t} + b\bar{b}$	2420	398	474	554	65.1	181	160	37.9	77
$t\bar{t}h$	818	151	189	125	24.4	59	37	8.42	16
Zh	329	84.6	87	37	3.68	22	73	16.8	32
Total bkg B	1.8×10^6	323064	259000	144000	17041	38800	116000	30237	49800
$\varsigma_{\text{stat}} = \sqrt{B}$	1330	568	509	380	131	197	341	174	223
$\varsigma_{\text{syst}} = 0.3\% \cdot B$	5290	969	778	433	51	116	349	91	149
$\varsigma_{\text{syst}} = 1\% \cdot B$	17600	3230	2590	1440	170	388	1160	302	498
$\varsigma_{\text{syst}} = 5\% \cdot B$	88000	16200	13000	7200	852	1940	5800	1510	2490
hh signal (κ_λ, κ_t)									
$(1, 1)_{\text{SM}}$	408	124	249	111	34.4	72	104	35.4	64
$(2, 1)$	194	58.9	127	67	25.7	47	77	24.5	44
$(5, 1)$	669	263	175	51	19.6	34	49	10.6	26
$(10, 1)$	5230	2000	1890	411	134	275	138	71.9	97
$(-5, 1)$	6210	2210	3050	847	361	595	381	102	181
$(1, 1.2)$	1010	316	626	270	83.6	179	216	89.1	139
$(1, 0.8)$	149	48.2	97	37	13.7	26	35	12.0	21

Table 4. Summary of signal region yields for the backgrounds and benchmark signals for the *baseline analysis* and *neural network analyses* (DNN) trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$. The yields shown are prior to binning in m_{hh} and normalised to integrated luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$. Dominant contributions to backgrounds are displayed together with the absolute size of the statistical ς_{stat} and benchmark systematic ς_{syst} uncertainties. Signals with benchmark couplings (κ_λ, κ_t) are shown, where $(1, 1)$ is the SM value.

λ_{hhh} constraints assuming $\kappa_t = 1$ and discusses the impact of systematic uncertainties. In subsection 4.3, we evaluate λ_{hhh}^i vs $\lambda_{hhh}^{j \neq i}$ discrimination power and discuss strategies for improvement. Finally, subsection 4.4 lifts the assumption of $\kappa_t = 1$ to examine the impact of top Yukawa y_t uncertainties on κ_λ constraints.

4.1 Signal and background rates

Table 4 shows the signal and background yields in the final signal region for the *baseline* (‘Baseline’ columns) and *neural network* (‘DNN’ columns) analyses prior to binning in m_{hh} . For the *baseline analysis* in the resolved category, we find the reducible $2b2j$ process comprises 49% of the background due to its formidable cross-section and only modest suppression from b -tagging. The irreducible $4b$ rate is comparable, comprising 45% of the background. These multijet processes constitute 94% of the backgrounds, similar to recent experimental analyses [41]. They also dominate in the intermediate and boosted categories

but interestingly, the $4b/2b2j$ ratio reduces to 0.25 and 0.21 respectively, suggesting light flavour jet systems have higher boost than those from b -jets. The top quark contribution, dominated by $t\bar{t}$, is around 3% for the resolved and intermediate categories. The single Higgs backgrounds contribute less than 1% of the total, with the most dominant contributions arising from the irreducible associated production $t\bar{t}h$ and the electroweak Zh processes.

For the *baseline analysis* assuming the SM $\kappa_\lambda = 1$ signal rate presented in table 4, we find a signal-to-background of $S/B = 2.3 \times 10^{-4}$ and significance of $S/\sqrt{B} = 0.30$ in the resolved category. For the intermediate category, the signal-to-background is higher at $S/B = 7.7 \times 10^{-4}$ but the significance $S/\sqrt{B} = 0.29$ is similar. For the boosted category, we find even higher signal-to-background $S/B = 9.0 \times 10^{-4}$ but again similar significance $S/\sqrt{B} = 0.31$. This pattern is consistent with the intermediate and boosted categories having higher signal purity due to better background rejection than the resolved but lower absolute yields. With $S/B \sim 10^{-3}$ to 10^{-4} and $B \gtrsim 10^4$ before binning in m_{hh} , we expect background statistical uncertainties $\zeta_{\text{stat}} = \sqrt{B}$ to be below the percent level and our analyses will be limited by systematic uncertainties.

For the *neural network analyses*, the DNN signal score p_{signal} is required to satisfy $p_{\text{signal}} > 0.75$. Two DNNs are considered: one trained on the SM $\kappa_\lambda = 1$ signal and the other on $\kappa_\lambda = 5$ where we expect the boundary of sensitivity to be for κ_λ limits. In table 4 for the resolved category, we find the $p_{\text{signal}} > 0.75$ requirement of the DNN trained on $\kappa_\lambda = 1$ gives a 61% SM $\kappa_\lambda = 1$ signal efficiency for 86% background rejection compared to the *baseline analysis*, providing a signal-to-background of $S_{\kappa_\lambda=1}/B = 9.6 \times 10^{-4}$ and significance of $S_{\kappa_\lambda=1}/\sqrt{B} = 0.49$. When a different signal hypothesis $\kappa_\lambda = 5$ is considered to that used in DNN training, we find the significance is lower $S_{\kappa_\lambda=5}/\sqrt{B} = 0.34$, as one would expect. The background rejection is dominated by the reduction in the $2b2j$ process such that the irreducible $4b$ component of the background now becomes dominant for the resolved category. For the DNN trained on $\kappa_\lambda = 5$, we find the $p_{\text{signal}} > 0.75$ requirement results in an 82% background rejection for 39% $\kappa_\lambda = 5$ signal efficiency, giving $S_{\kappa_\lambda=5}/B = 8.1 \times 10^{-4}$ and $S_{\kappa_\lambda=5}/\sqrt{B} = 0.46$. As one would expect, this is higher for the $\kappa_\lambda = 5$ signal than using the DNN trained on $\kappa_\lambda = 1$, but this is only the case for the resolved category. Interestingly for the intermediate and boosted categories, the DNN trained on $\kappa_\lambda = 1$ results in higher signal yields for all signal hypotheses listed in table 4. Such categories have a greater signal acceptance times efficiency for $\kappa_\lambda = 1$ than for $\kappa_\lambda = 5$ scenarios (figure 6) as the former κ_λ case has a harder Higgs p_T spectrum (figure 3).

Turning to the yields binned in m_{hh} , table 5 compares these across the three analyses and three categories for the $\kappa_\lambda = 1, 2, 5$ signals. For the DNN trained on $\kappa_\lambda = 1$, the yields are suppressed in the low m_{hh} bins such that there is no signal in $m_{hh} \in [200, 250]$ GeV. This is expected given the background rate is high but the signal rate is low in these bins. Therefore, only training the DNN on $\kappa_\lambda = 1$ signals is suboptimal for signals that occupy lower m_{hh} values such as $\kappa_\lambda = 5$. We see that the $\kappa_\lambda = 5$ training retains the $\kappa_\lambda = 5$ signal in the first $m_{hh} \in [200, 250]$ bin, which is important for κ_λ constraints that will be further discussed in the next subsections.

Analysis	Baseline				DNN trained on $\kappa_\lambda = 5$				DNN trained on $\kappa_\lambda = 1$			
m_{hh} bin [GeV]	B	$S_{\kappa_\lambda=1}$	$S_{\kappa_\lambda=2}$	$S_{\kappa_\lambda=5}$	B	$S_{\kappa_\lambda=1}$	$S_{\kappa_\lambda=2}$	$S_{\kappa_\lambda=5}$	B	$S_{\kappa_\lambda=1}$	$S_{\kappa_\lambda=2}$	$S_{\kappa_\lambda=5}$
Resolved												
[200, 250]	369000	3.1	3.6	133	100000	0.76	0.662	50.9	98	0.0	0.0	0.0
[250, 300]	687000	28.4	7.9	231	112000	10.0	2.0	83.9	5660	4.1	0.7	11.2
[300, 350]	310000	68.6	21.6	137	42500	23.6	9.6	69.1	85200	25.6	10.2	55.4
[350, 400]	154000	87.8	42.7	66.5	25300	32.5	16.6	25.2	75200	59.7	27.5	33.5
[400, 500]	179000	170	88.3	82.7	18200	52.1	25.5	31.4	81900	128.3	67.7	62.2
[> 500]	37600	50.1	30.0	17.3	1310	5.2	4.6	2.5	11400	30.9	21.0	12.5
Intermediate												
[200, 500]	36600	28.7	14.1	14.9	5090	8.5	6.0	5.53	8250	16.7	9.0	8.9
[500, 600]	69300	56.4	35.3	24.2	8310	20.0	14.2	9.2	20200	37.6	25.0	16.0
[> 600]	38600	26.0	17.4	12.1	3640	5.9	5.5	4.89	10300	18.0	13.3	9.1
Boosted												
[500, 800]	69400	74.3	49.8	27.9	25800	30.8	20.4	8.5	40500	52.5	34.7	17.6
[> 800]	46600	29.1	27.3	21.2	4430	4.6	4.1	2.15	9300	11.0	8.9	7.9

Table 5. Summary of signal region event yields binned by m_{hh} for the total background B and benchmark signals for different κ_λ couplings $S_{\kappa_\lambda=i}$. These are displayed for the *baseline* and *neural network* (DNN) trained on $\kappa_\lambda = 5$, and trained on $\kappa_\lambda = 1$ for comparison. The yields are normalised to an integrated luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$ for the three categories.

We now evaluate the binned statistical significances defined by $Z = S/\sqrt{B + (\zeta_b B)^2}$ including background systematics ζ_b and display these in table 7. For each category, the lowest two rows show the significances combined in two ways for comparison:

- *Quadrature sum.* The significance for each m_{hh} bin i is summed in quadrature to give $Z = \sqrt{\sum_i Z_i^2}$. This accounts for m_{hh} shape information, where for simplicity we assume no correlations in systematics.
- *One bin.* The yields in each bin i are summed $S = \sum_i S_i, B = \sum_i B_i$ before the significance Z is calculated such that no m_{hh} shape information is considered.

Finally, the bottom three rows of table 7 combine the significances across all three categories. The quadrature sum considers all the m_{hh} bins in all three categories, while the combined one bin approach is the quadrature sum of the *one bin* significances from each category. For simplicity, we neglect statistical and systematic uncertainties on the signal in the significances as analyses are in background-dominated $S \ll B$ regimes. Given $B \gg 1$, we assume statistical uncertainties follow the Gaussian approximation $\zeta_{\text{stat}}^B = \sqrt{B}$.

For the systematic uncertainties, we adopt a simple approach that assumes the nominal benchmark values ζ_b summarised in table 6. We assume $\zeta_b = 0.3\%$ for the resolved category, which is very demanding but in accord with ref. [57]; this is a HL-LHC extrapolation of the recent ATLAS analysis which reports 2% systematics [41]. For the boosted category, a similar ATLAS analysis [41] currently reports up to 18% systematics, dominated by the data-driven background estimate method and jet mass resolution [41], while the current CMS analysis report up to 7% systematics [43]. For our work, we assume that this can be controlled to the level of $\zeta_b = 5\%$ at the HL-LHC. For the intermediate

Category	Systematic ζ_b
Resolved	0.3%
Intermediate	1%
Boosted	5%

Table 6. The nominal relative systematic uncertainties on the background estimate ζ_b assumed when performing the χ^2 analysis for the different analysis categories.

category, we assume a benchmark value between what we apply for resolved and boosted of $\zeta_b = 1\%$. In experimental implementations [39–45], the multijet background is typically modelled via data-driven methods so is not subject to theoretical uncertainties such as renormalisation scale, parton shower and parton distribution functions. Nonetheless, data-driven procedures have systematics due to unknown composition and extrapolation over correlated variables.

For the *baseline analysis* significances in table 7, the total significance is higher for the quadrature sum approach compared to the one bin counterpart, demonstrating the importance of m_{hh} shape information for sensitivity. With $\zeta_b = 0\%$, the $Z_{\kappa_\lambda=1}$ significance evaluated using the $\kappa_\lambda = 1$ signal for the resolved category is 0.54 compared to 0.37 with the nominal 0.3% systematics and just 0.08 without binning in m_{hh} . The intermediate and boosted categories have comparable significances of 0.29 and 0.31 respectively with $\zeta_b = 0\%$, but dilutes to 0.12 and 0.02 with their nominal systematics of table 6. The significance combining all categories is 0.69 for $\zeta_b = 0\%$ and 0.39 with nominal systematics.

Moving to the *neural network analyses*, we find the significances are higher than those of the corresponding *baseline analysis* across all m_{hh} bins. This is especially important when systematics are included, where we see the resolved category $Z_{\kappa_\lambda=1}$ significance improve from 0.54 to 0.59 for the DNN trained on $\kappa_\lambda = 1$ with 0% systematics but the improvement is 0.37 to 0.48 with the nominal 0.3% systematics. When systematics are included, this suggests that the background rejection is more important to improve significance. We find that the DNN trained on $\kappa_\lambda = 5$ provides higher significances for $\kappa_\lambda = 5$ signals ($Z_{\kappa_\lambda=5} = 0.44$) than the DNN trained on $\kappa_\lambda = 1$ ($Z_{\kappa_\lambda=5} = 0.30$) in the resolved category. The intermediate category has significances that are numerically about half those of the resolved category: $Z_{\kappa_\lambda=1} = 0.24$ (vs 0.48) for DNN trained on $\kappa_\lambda = 1$ intermediate (vs resolved). With all categories combined (bottom rows of table 7), the DNN trained on $\kappa_\lambda = 1$ gives $Z_{\kappa_\lambda=1} = 0.75$ with 0% systematics and is diluted to 0.53 with systematics, which improves upon the *baseline analysis* significance of 0.39. The significance without binning in m_{hh} is lower at 0.32, suggesting that additional discrimination power remained in the m_{hh} shape not captured by a cut on the DNN score.

4.2 χ^2 analysis of λ_{hhh} fixing $\kappa_t = 1$

We now perform a χ^2 analysis to evaluate projections of Higgs self-coupling λ_{hhh} constraints. We adapt the definition in eq. (3.1) into one for benchmarking our analyses:

$$\chi^2 = \sum_i \left[\frac{(S - S_{\text{SM}})^2}{S + B + (\zeta_b B)^2 + (\zeta_s S)^2} \right]_i. \quad (4.1)$$

Analysis Trained on	Baseline —		DNN $\kappa_\lambda = 1$		DNN $\kappa_\lambda = 5$	
Resolved						
m_{hh} bin [GeV]	$Z_{\kappa_\lambda=1}$	$Z_{\kappa_\lambda=5}$	$Z_{\kappa_\lambda=1}$	$Z_{\kappa_\lambda=5}$	$Z_{\kappa_\lambda=1}$	$Z_{\kappa_\lambda=5}$
[200, 250]	0.00	0.11	0.00	0.00	0.00	0.12
[250, 300]	0.01	0.10	0.05	0.15	0.02	0.18
[300, 350]	0.06	0.13	0.07	0.14	0.10	0.29
[350, 400]	0.14	0.11	0.17	0.09	0.18	0.14
[400, 500]	0.25	0.12	0.34	0.17	0.36	0.22
> 500	0.22	0.08	0.28	0.11	0.14	0.07
Quadrature sum m_{hh} bins (0% syst)	0.54	0.51	0.59	0.37	0.47	0.53
Quadrature sum m_{hh} bins	0.37	0.27	0.48	0.30	0.44	0.44
One bin $m_{hh} > 200$ GeV	0.08	0.12	0.27	0.19	0.12	0.25
Intermediate						
[200, 500]	0.07	0.04	0.14	0.07	0.10	0.06
[500, 600]	0.08	0.03	0.15	0.06	0.16	0.07
> 600	0.06	0.03	0.12	0.06	0.08	0.07
Quadrature sum m_{hh} bins (0% syst)	0.29	0.14	0.37	0.17	0.27	0.15
Quadrature sum m_{hh} bins	0.12	0.06	0.24	0.12	0.21	0.12
One bin $m_{hh} > 200$ GeV	0.07	0.03	0.17	0.08	0.16	0.09
Boosted						
[500, 800]	0.02	0.01	0.03	0.01	0.02	0.01
> 800	0.01	0.01	0.02	0.02	0.02	0.01
Quadrature sum m_{hh} bins (0% syst)	0.31	0.14	0.28	0.12	0.20	0.062
Quadrature sum m_{hh} bins	0.02	0.01	0.03	0.02	0.03	0.01
One bin $m_{hh} > 500$ GeV	0.02	0.01	0.03	0.01	0.02	0.01
Combined quadrature sum (0% syst)	0.69	0.55	0.75	0.42	0.58	0.56
Combined quadrature sum	0.39	0.27	0.53	0.32	0.49	0.46
Combined one bin	0.11	0.13	0.32	0.20	0.20	0.27

Table 7. Statistical significances defined by $Z = S/\sqrt{B + (\zeta_b B)^2}$ for $\mathcal{L} = 3000 \text{ fb}^{-1}$ using the signal and background counts of table 5. Here, Z_i corresponds to using the $\kappa_\lambda = i$ signal samples and the nominal systematics ζ_b are taken from table 6. These are evaluated for the baseline analysis, the DNN analysis where the DNN was trained with the $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$ samples. The results are shown for the resolved, the intermediate and boosted categories which are separated by the corresponding headings. The lowest two rows of each category presents the significances combined by the ‘quadrature sum’ of the significances binned in m_{hh} , and second ‘one bin’ sums the yields into one inclusive bin before calculating Z . The bottom three rows shows significances combined across all three categories.

Here, B is the total background rate, S_{SM} is the SM signal rate, and S is the signal rate we wish to distinguish from the SM counterpart. These are evaluated in each m_{hh} bin i using the yields in table 5. Like the previous significance calculations of subsection 4.1, we neglect statistical and systematic uncertainties on the signal as we are in background-dominated $S \ll B$ regimes. For simplicity, we perform a statistical combination by summing the χ^2 values assuming constant and uncorrelated uncertainties across bins $\{i\}$. We first combine

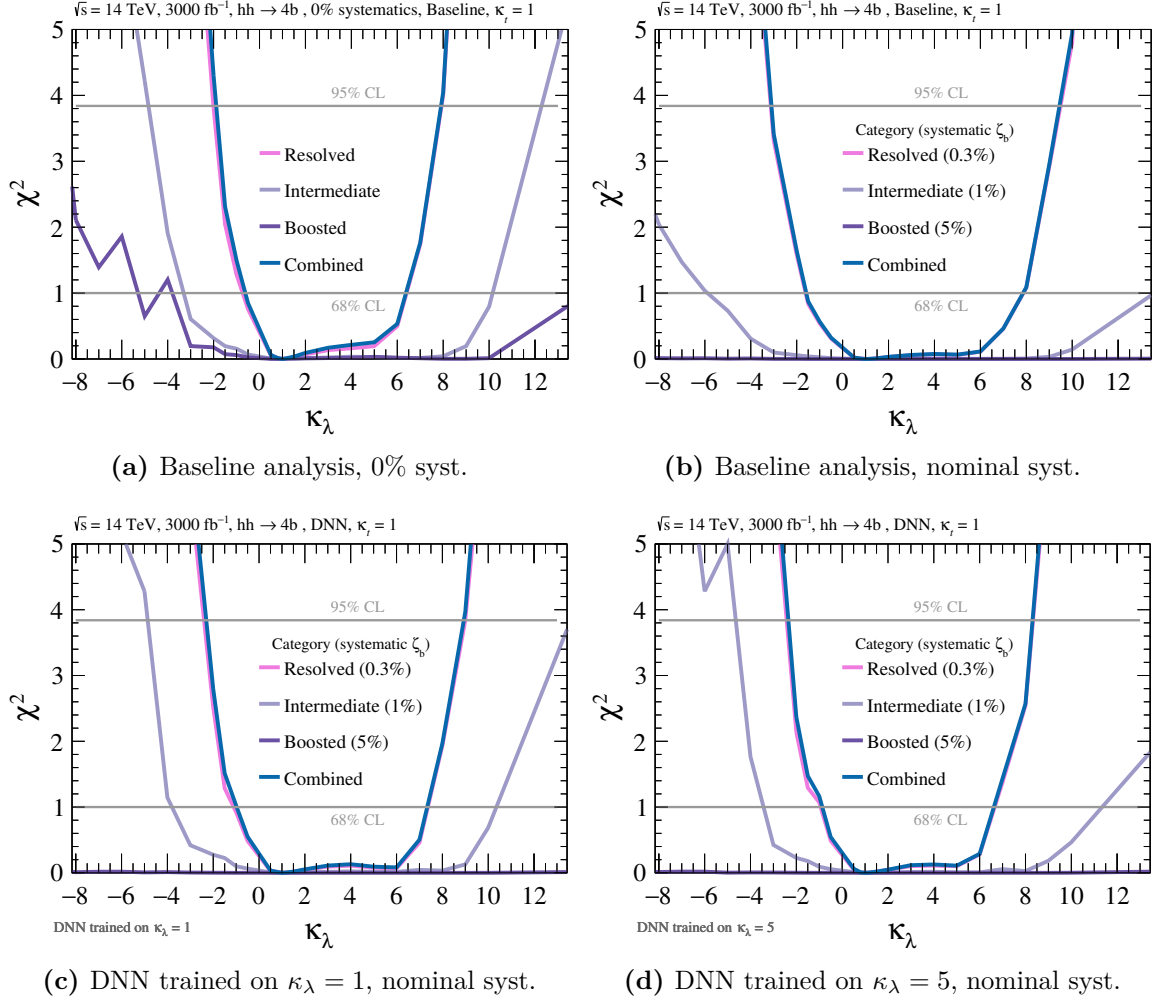


Figure 11. The χ^2 vs κ_λ distributions with fixed $\kappa_t = 1$ for different analyses and assumed systematics using $\mathcal{L} = 3000 \text{ fb}^{-1}$. This is displayed in the resolved (pink), intermediate (lilac), boosted (purple) categories, together with their combination (blue). The *baseline analysis* is shown assuming (a) 0% background systematics and (b) the nominal systematics of 0.3%, 1% and 5% for the resolved, intermediate and boosted categories, respectively (table 6). The *neural network analysis* is shown with the nominal systematics for the DNN trained on (c) $\kappa_\lambda = 1$ and (d) $\kappa_\lambda = 5$. The grey horizontal lines at $\chi^2 = 1, 3.84$ indicate the 68% CL, 95% CL thresholds respectively. The unphysical spikes for $\kappa_\lambda < 4$ for the boosted χ^2 line in subfigure (a) and $\kappa_\lambda < -5$ for subfigure (d) are due to limited MC statistics.

the yields in the m_{hh} bins shown in table 5 separately for the resolved, intermediate and boosted categories to explore their complementary sensitivity. Then we combine the three categories to obtain the final combined constraint.

For the *baseline analysis*, figure 11a shows the χ^2 vs κ_λ distribution without systematics. The resolved category is most constraining $\kappa_\lambda \in [-0.75, 6.5]$ at 68% CL, followed by the intermediate $\kappa_\lambda \in [-3.3, 10.2]$ at 68% CL with boosted being the weakest category. To gain intuition for the sensitivity between the three categories, we note that figure 3 shows

that signals with $\kappa_\lambda \gtrsim 5$ have the lowest Higgs p_T spectrum. Therefore, the intermediate and boosted scenarios have lower acceptance for these scenarios where the upper boundary of κ_λ sensitivity is expected. Meanwhile, for negative κ_λ , we see that the Higgs bosons in $\kappa_\lambda \lesssim -1$ scenarios have modest boost (figure 3). Therefore, the acceptance is higher for all categories and the constraints are stronger for negative κ_λ . Although scenarios near the SM value $\kappa_\lambda = 1$ have the most boosted Higgs bosons and signal acceptance is high (figure 6), these κ_λ values are very challenging to constrain in a χ^2 analysis because the total signal cross-sections are small and change slowly with respect to κ_λ due to destructive interference (subsection 2.1). Inspecting table 4, we find factors of two change in signal yield between the challenging hypotheses of $\kappa_\lambda = 1, 2, 5$. However, the yields binned in m_{hh} (table 5) suggest that our shape analysis does have modest sensitivity to signal distribution changes for these κ_λ scenarios even if it is insufficient for 68% CL sensitivity.

Figure 11b shows the χ^2 vs κ_λ distributions assuming the nominal systematics of table 6. The 0.3% systematics for the resolved category dilutes the constraints to $\kappa_\lambda \in [1.6, 7.9]$ at 68% CL. Meanwhile, the 1% systematics for the intermediate category weakens the constraint to $\kappa_\lambda \in [-5.8, 13.6]$ at 68% CL. However, the boosted category loses all constraining power in the considered κ_λ range for the assumed 5% systematics. Overall for the *baseline analysis*, the statistical combination of all three categories leads to a 68% CL constraint of $\kappa_\lambda \in [-0.6, 6.4]$ for no systematics. This degrades to $\kappa_\lambda \in [-1.6, 7.9]$ when assuming the nominal systematics, driven by the sensitivity of the resolved category.

Turning to the *neural network analysis*, we evaluate the χ^2 vs κ_λ distribution assuming the nominal systematics with the DNN trained on $\kappa_\lambda = 1$ (figure 11c) and $\kappa_\lambda = 5$ (figure 11d). Statistically combining all three categories achieves 68% CL constraints of $\kappa_\lambda \in [-1.0, 7.4]$ and $\kappa_\lambda \in [-0.8, 6.6]$ for the DNN trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$ respectively. These improve upon the *baseline analysis* of $\kappa_\lambda \in [-1.6, 7.8]$, with the $\kappa_\lambda = 5$ DNN training surpassing those achieved by the $\kappa_\lambda = 1$ DNN training. This shows that dedicated optimisation for a signal closer to the boundary of sensitivity can improve κ_λ constraints. We choose to train on the $\kappa_\lambda = 5$ signal due to its substantial triangle diagram contribution and therefore has a significant fraction of events at low m_{hh} . The kinematics are also reasonably representative of higher κ_λ values (figure 3b). Training on $\kappa_\lambda = 5$ instead of $\kappa_\lambda = 1$ encourages the neural network to retain low m_{hh} events, which has larger shape differences between different κ_λ values. This underscores the importance of optimisation away from the signals with SM couplings for κ_λ constraints in the $hh \rightarrow 4b$ final state.

A future improvement may consider designing multiple independent neural networks, each optimised for one sampled value of κ_λ . A more sophisticated approach may extend the idea of a parameterised neural network [80] to construct an observable that is a well-behaved function of κ_λ and provides good signal discrimination power, but this is beyond the scope of this work. While we find that sensitivity is driven by the resolved category, it is noteworthy that the DNN substantially improves the κ_λ constraints of the intermediate category from $\kappa_\lambda \in [-5.8, 13.6]$ of the *baseline analysis* to $\kappa_\lambda \in [-3.8, 10.4]$ for the DNN trained on $\kappa_\lambda = 1$ and $\kappa_\lambda \in [-3.4, 11.3]$ for the DNN trained on $\kappa_\lambda = 5$. In a wider context, we note that intermediate category has close to comparable λ_{hhh} constraints with projected HL-LHC constraints for the di- b -quark plus diboson $hh \rightarrow bbVV$ final states [58].

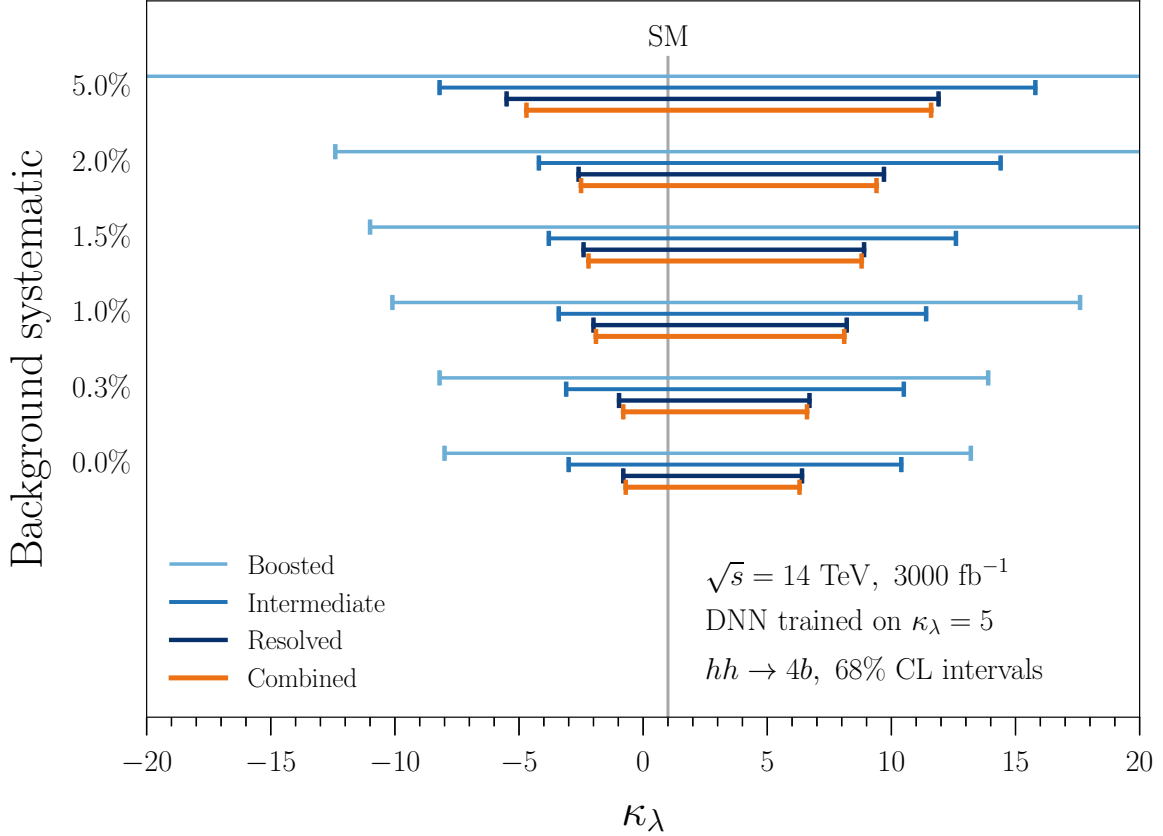


Figure 12. Summary of 68% CL intervals ($\chi^2 < 1$) on κ_λ with fixed $\kappa_t = 1$ for different assumed background systematics from 0% to 5% for the *neural network analysis* trained on the $\kappa_\lambda = 5$ signal. This is shown for the resolved (dark blue), intermediate (medium blue), boosted (light blue) categories along with their combination (orange). The luminosity is assumed to be 3000 fb^{-1} . Lines without endcaps mean the 68% CL limit is outside the considered range of $\kappa_\lambda \in [-20, 20]$. The vertical grey line denotes the SM value of κ_λ . The nominal systematics assumed in this work are shown in table 6. In appendix C, figure 27 shows the version of this plot with the DNN trained on the $\kappa_\lambda = 1$ signal.

While we have assumed the nominal systematics of table 6, figure 12 now evaluates the 68% CL κ_λ limits for different systematics ranging from 0% to 5% with the DNN trained on $\kappa_\lambda = 5$. We find that the constraints for the resolved category degrade to $\kappa_\lambda \in [-5.5, 11.9]$ for $\zeta_b = 5\%$ and the intermediate category constraints weaken to $\kappa_\lambda \in [-8.2, 15.8]$. This shows how important the assumed level of systematics are for constraining λ_{hhh} due to small $S/B \simeq 10^{-3}$. Meanwhile, the boosted category has no constraining power within $\kappa_\lambda \in [-20, 20]$ when systematic uncertainties are assumed at the nominal 5% level. If these could be improved to the 2% level, the boosted category can constrain $\kappa_\lambda > -12.4$. We also find that if the systematics are controlled to 1%, there is additionally sensitivity to positive $\kappa_\lambda < 17.6$. The difficulty in forecasting experimental systematics at the HL-LHC leads to significant uncertainty for projections of κ_λ constraints. One can also interpret these systematics as benchmark targets for future analyses to improve λ_{hhh} constraints.

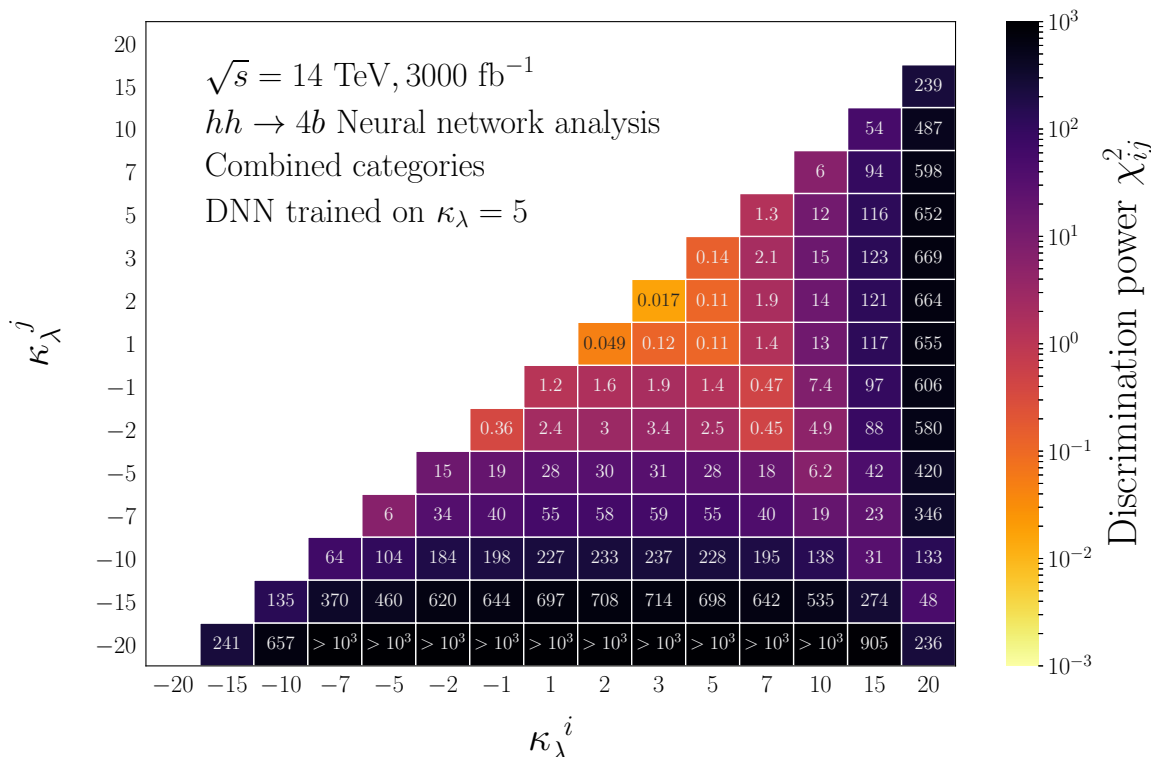


Figure 13. Discrimination power χ^2_{ij} (higher is better) between two self-coupling hypotheses λ^i_{hhh} vs λ^j_{hhh} for the *neural network analysis* after combining the three categories. The DNN is trained on the $\kappa_\lambda = 5$ sample, where the signal score is required to be $p_{\text{signal}} > 0.75$. The χ^2_{ij} is calculated from eq. (3.1) assuming $\kappa_t = 1$ and background systematics of $\zeta_b = 0.3\%, 1\%$ and 5% for the resolved, intermediate and boosted categories, respectively. For comparison, this figure for the *baseline* and *neural network* analyses trained on $\kappa_\lambda = 1$, are found in appendix C.

4.3 λ^i_{hhh} vs $\lambda^{j \neq i}_{hhh}$ discrimination power

So far, our χ^2 distributions (figure 11) have been evaluated with respect to the SM $\kappa_\lambda = 1$ signal i.e. assuming we observe the SM signal in data. We now generalise this to evaluate the discrimination power between any two κ_λ hypotheses, which we present in figure 13. While we train our DNN analysis on one class of signal (here $\kappa_\lambda = 5$), it is important to quantify the signal characterisation capability should λ_{hhh} deviate from the SM or the trained signal scenario. We calculate the χ^2_{ij} between two signal hypotheses (i, j) shown on the axes of figure 13. We present this for the *neural network analysis* trained on $\kappa_\lambda = 5$, combining the three categories and assuming the nominal systematics in table 6. Larger values of χ^2_{ij} indicate that it is easier to discriminate between two hypotheses (i, j), while χ^2_{ij} values around unity or below indicate difficulty in discriminating.

Overall, we find figure 13 shows good discrimination power where $\chi^2_{ij} \gg 1$ for scenarios with $\kappa_\lambda \geq 10$ and $\kappa_\lambda \leq -5$, whose sensitivity is driven by the total cross-section. This corresponds to κ_λ values far away from destructive interference and the cross-section changing rapidly with respect to κ_λ . We also note small $\chi^2_{ij} < 1$ occurs around $\kappa_\lambda = -2$ and $\kappa_\lambda = 7$, which has $\chi^2_{ij} \simeq 0.5$. This arises physically due to the total cross-sections being

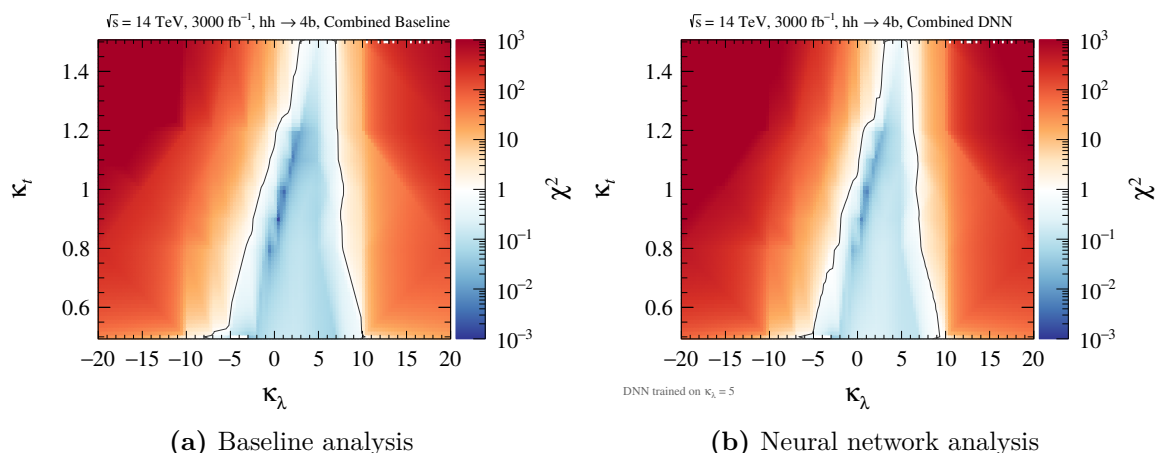


Figure 14. The combined χ^2 distributions for the (a) *baseline* and (b) *neural network* analyses in the κ_t vs κ_λ plane, assuming $\mathcal{L} = 3000 \text{ fb}^{-1}$. The χ^2 statistically combines the resolved, intermediate and boosted categories with 0.3%, 1% and 5% background systematic uncertainties, respectively. The DNN is trained on $\kappa_\lambda = 5$. The grey contour indicates $\chi^2 = 1$ corresponding to 68% CL. A breakdown by analysis category is displayed in appendix C.

similar $\sigma(\kappa_\lambda = -2)/\sigma(\kappa_\lambda = 7) \simeq 0.84$ combined with acceptance being higher for negative κ_λ (figure 6) due to harder Higgs p_T spectra (figure 3). These sign degeneracies are not symmetric about zero due to maximal destructive interference being near $\kappa_\lambda \sim 2.5$. Generally, scenarios that are difficult to distinguish fall along a northwest-to-southeast corridor of low χ^2_{ij} values. There is also a cluster around $1 \leq \kappa_\lambda < 4$ with $\chi^2_{ij} < 1$ arising from both total cross-sections and changes with respect to κ_λ being suppressed due to near-maximal destructive interference (subsection 2.1).

For future work, it would be interesting to exploit m_{hh} shape information better via dedicated optimisations to break degeneracies in the κ_λ parameter space further. Moreover, one could imagine a DNN trained just on a single κ_λ hypothesis $\kappa_\lambda = 5$ is suboptimal for discriminating against two signal hypotheses unrelated to $\kappa_\lambda = 5$. It would be desirable to design analyses, such as generalising the current neural network strategy, to optimise the discrimination power across a wider variety of κ_λ values.

4.4 κ_λ constraints allowing $\kappa_t \neq 1$

Given the rapid scaling of di-Higgs cross-sections with the top Yukawa coupling y_t shown by eq. (2.2), we expect modifications in y_t to impact the projected λ_{hhh} constraints, which we now explore. Figure 14a shows the χ^2 distribution in the κ_t vs κ_λ plane for the *baseline analysis* after combining the resolved, intermediate and boosted categories, assuming the nominal systematics of table 6. The contours are displayed after linear interpolation of the points sampled from this parameter space with the density shown in figure 2. The grey line marks out the $\chi^2 = 1$ contour that corresponds to 68% CL. Consistent with the one-dimensional distributions (figure 11b), we see the red colours representing regions that have high χ^2 values and are statistically excluded, while the blue shades indicate parameters that have low sensitivity. The dark blue strip of very low χ^2 that runs from

Analysis	DNN trained on	κ_t	68% CL on κ_λ			
			Resolved	Intermediate	Boosted	Combined
Baseline	—	1.2	[0.2, 7.2]	[−3.7, 11.3]	—	[0.2, 7.2]
DNN	$\kappa_\lambda = 1$	1.2	[0.5, 7.1]	[−1.2, 10.3]	—	[0.7, 7.1]
DNN	$\kappa_\lambda = 5$	1.2	[0.4, 6.3]	[−3.1, 10.4]	—	[0.6, 6.3]
DNN 0% syst	$\kappa_\lambda = 5$	1.2	[0.6, 5.8]	[−0.8, 10.3]	[−3.7, 11.0]	[0.9, 5.8]
Baseline	—	1.0	[−1.6, 7.9]	[−5.8, 13.6]	—	[−1.6, 7.9]
DNN	$\kappa_\lambda = 1$	1.0	[−1.1, 7.4]	[−3.8, 10.4]	—	[−1.0, 7.4]
DNN	$\kappa_\lambda = 5$	1.0	[−1.0, 6.6]	[−3.4, 11.3]	—	[−0.8, 6.6]
DNN 0% syst	$\kappa_\lambda = 5$	1.0	[−0.8, 6.4]	[−3.0, 10.4]	[−8.0, 13.2]	[−0.7, 6.3]
Baseline	—	0.8	[−3.3, 7.8]	[−10.0, 11.9]	—	[−3.3, 7.8]
DNN	$\kappa_\lambda = 1$	0.8	[−2.7, 7.5]	[−5.2, 10.4]	—	[−2.6, 7.4]
DNN	$\kappa_\lambda = 5$	0.8	[−2.8, 7.2]	[−5.2, 10.4]	—	[−2.6, 7.2]
DNN 0% syst	$\kappa_\lambda = 5$	0.8	[−2.3, 7.0]	[−4.3, 10.3]	[−8.4, 13.2]	[−2.2, 6.9]

Table 8. Summary of 68% CL limits on κ_λ for *baseline analysis* and *neural network analysis* (DNN) by separate categories and combined. These are shown for $\kappa_t = 1.2, 1.0, 0.8$ to demonstrate the impact of uncertainties on the top Yukawa. Systematic uncertainties on the background of 0.3%, 1% and 5% are assumed for the resolved, intermediate and boosted categories, respectively. The DNN analysis trained on $\kappa_\lambda = 5$ is also shown with 0% systematics for comparison. A ‘—’ in the 68% CL columns indicates no constraint in the range of κ_λ considered, $\kappa_\lambda \in [-20, 20]$.

around $(\kappa_\lambda, \kappa_t) \sim (0, 0.8)$ to $(3, 1.2)$ indicates a region that is very similar to the SM $(1, 1)$ parameters, given this is our reference value for calculating the χ^2 .

Of particular interest is the shape of the $\chi^2 = 1$ contour in the two-dimensional κ_t vs κ_λ plane: as κ_t increases, the 68% CL constraint on κ_λ improves, indicated by the region satisfying $\chi^2 < 1$ decreasing. To make quantitative comparisons, we choose benchmark κ_t values of 0.8 and 1.2 based on current uncertainties from direct $pp \rightarrow t\bar{t}h$ measurements of y_t being on the order of 20% [67, 69]. We note that while measurements of gluon fusion production have smaller uncertainties, these indirectly probe y_t via the one-loop gluon-gluon-Higgs interaction.

For the *baseline analysis*, the combined 68% CL constraint goes from $\kappa_\lambda(\kappa_t = 1) \in [-1.6, 7.9]$ to $\kappa_\lambda(\kappa_t = 1.2) \in [0.2, 7.2]$ and $\kappa_\lambda(\kappa_t = 0.8) \in [-3.3, 7.8]$. These constraints are also summarised in table 8 to facilitate comparison. We note that the change in sensitivity is more significant for negative λ_{hh} , where constraining the sign of κ_λ to be positive is possible if $\kappa_t = 1.2$. To provide a clearer intuition of the impact, we can define how the absolute width $\delta\kappa_\lambda = |\kappa_\lambda^+ - \kappa_\lambda^-|$ of the 68% CL intervals change, where $\kappa_\lambda^{+(-)}$ is the upper (lower) bound on κ_λ . With this metric, we find that the absolute width of the 68% CL constraint $\delta\kappa_\lambda(\kappa_t = 1) = 9.5$ decreases by 26% for $\kappa_t = 1.2$ and increases by 17% for $\kappa_t = 0.8$. The results of the *neural network analysis* are shown in figure 14b, where the DNN trained on the $\kappa_\lambda = 5$ signal is used. As table 8 also summarises, we find the 68% CL constraint goes from $\kappa_\lambda(\kappa_t = 1) \in [-0.8, 6.6]$ to $\kappa_\lambda(\kappa_t = 1.2) \in [0.6, 6.3]$ and

$\kappa_\lambda(\kappa_t = 0.8) \in [-2.6, 7.2]$. Using the $\delta\kappa_\lambda$ metric, we find $\delta\kappa_\lambda(\kappa_t = 1) = 7.4$ decreases by 23% for $\kappa_t = 1.2$ and increases by 32% for $\kappa_t = 0.8$. These changes in $\delta\kappa_\lambda$ are non-trivial to interpret but they suggest that the size of the systematic on κ_λ projections due to uncertainties on y_t impacting the signal modelling can be substantial.

A further observation in figure 14a is that we can extract a 68% CL upper limit on the top Yukawa $\kappa_t < 1.25$ assuming fixed $\kappa_\lambda = 1$. This improves slightly to $\kappa_t < 1.22$ using the DNN trained on $\kappa_\lambda = 5$. It is interesting that this additional result of our work arises without considering dedicated optimisation for constraining κ_t . That the di-Higgs process has constraining power to put a better than 25% upper bound on κ_t is helped by the fast quartic dependence of the box diagram contribution to the cross-section $\sigma_{\text{box}} \propto y_t^4$. The ability to constrain κ_t depends strongly on the value of κ_λ : for example, we find no κ_t constraint if $\kappa_\lambda = 5$. Although our $hh \rightarrow 4b$ analysis is not expected to compete with other higher statistics channels in constraining y_t (directly via $t\bar{t}h$ or indirectly via gluon fusion), it provides an additional independent probe that could contribute in a global fit.

Now we extend our discussion by using figure 15, which overlays two-dimensional 68% CL contours to compare the *baseline analysis* (dashed lines) with DNN analysis trained on $\kappa_\lambda = 5$ (solid lines). This figure also separates the contours by category, but as the boosted has little constraining power (in part due to the conservative 5% systematic), only comparisons between resolved and intermediate are of interest. Similar to the one-dimensional χ^2 distributions, sensitivity is driven by the resolved category across $\kappa_t \neq 1$ with the intermediate category contribution being non-negligible for negative κ_λ .

Assuming $\kappa_t = 1$, the resolved category *baseline analysis* is $\kappa_\lambda \in [-1.6, 7.9]$ and improves by 20% (in $\delta\kappa_\lambda$) to $\kappa_\lambda \in [-1.0, 6.6]$ for the DNN trained on $\kappa_\lambda = 5$, which also is summarised in table 8. While the intermediate category has comparably weaker constraints of $\kappa_\lambda \in [-5.8, 13.6]$ for the *baseline analysis*, the DNN trained on $\kappa_\lambda = 5$ improves this by 24% to $\kappa_\lambda \in [-3.4, 11.3]$. This improvement is largely uniform for different κ_t . For negative κ_λ , the slope of the intermediate contour in the κ_t vs κ_λ plane is similar to that of the resolved. This suggests that the two categories are sensitive to similar effects of the constructive interference occurring for negative κ_λ . By contrast, the fact that the intermediate contour is nearer vertical for positive κ_λ shows it is less sensitive to changes in κ_t even if the constraints are weaker than resolved. Additionally, we see that it is the resolved category that drives the $\kappa_t < 1.22$ 68% CL constraint for $\kappa_\lambda = 1$.

5 Conclusion

This paper presented a comprehensive analysis of Higgs pair production in $4b$ final states to evaluate and improve Higgs self-coupling (λ_{hhh}) constraints at the HL-LHC. We extended current strategies in several directions and now summarise our conclusions.

We performed a detailed comparison of λ_{hhh} constraints across event categories that have zero (resolved), exactly one (intermediate) and two or more (boosted) boosted Higgs bosons in the $4b$ channel for the first time. The resolved analysis was found to be most constraining, followed by the intermediate and then the boosted. We used deep neural networks for signal-background separation and achieved an 86% background rejection

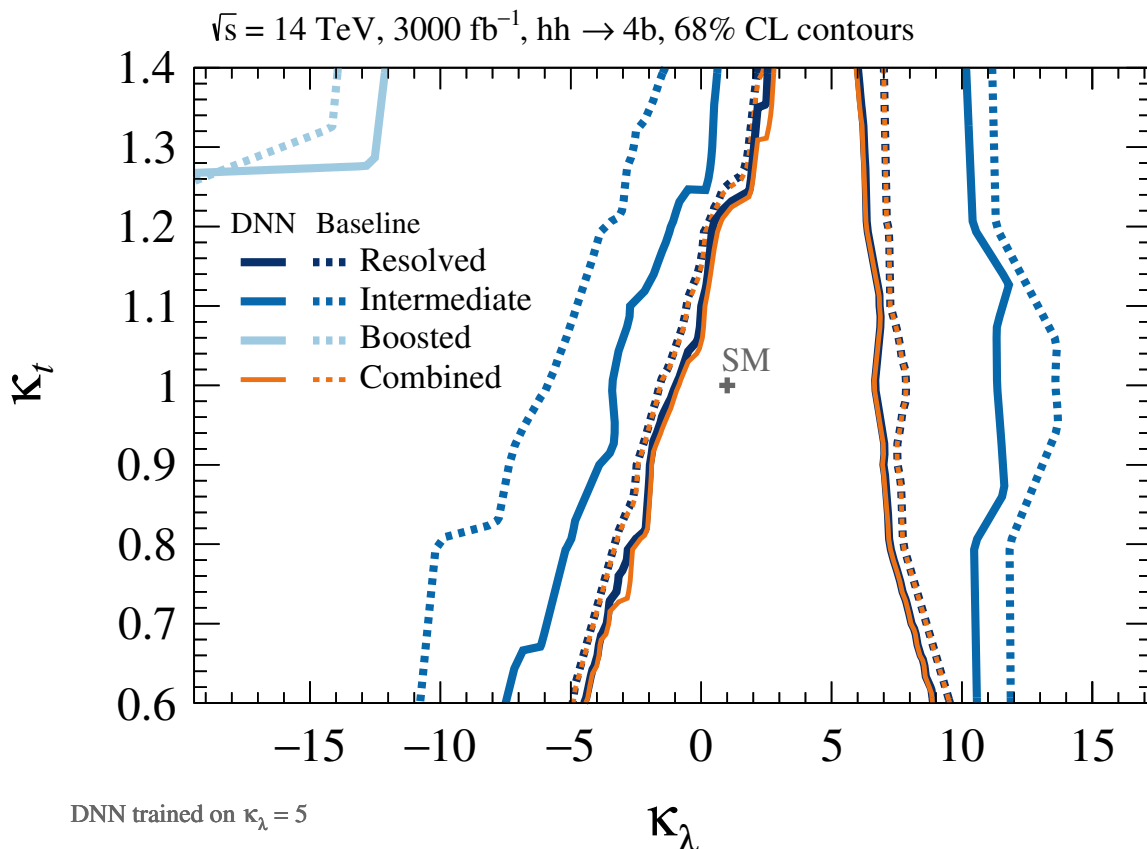


Figure 15. Summary of 68% CL ($\chi^2 = 1$) contours in the κ_t vs κ_λ plane. These are displayed for the resolved (dark blue), intermediate (medium blue), boosted (light blue) categories and their combination (orange) for the *baseline analysis* (dashed) and *neural network analysis* (DNN) trained on $\kappa_\lambda = 5$ (solid). A luminosity of 3000 fb^{-1} is assumed along with systematic uncertainties of 0.3%, 1% and 5% for the resolved, intermediate and boosted categories, respectively. The cross indicates the SM prediction. The combined χ^2 distributions are displayed in figure 14 and the χ^2 distributions for each analysis category are shown in appendix C.

for 61% SM signal efficiency in the resolved category compared to the *baseline analysis*. This improved the signal significance from 0.69 (0.39) to 0.75 (0.53) without (with) the nominal systematics assumed, which can contribute in combination with other channels. Importantly, analyses optimised for discovery of SM hh production can be suboptimal for constraining λ_{hhh} in the $4b$ channel. We improved constraints by optimising on signals with λ_{hhh} values that are five times the SM and have lower boosts than those closer to SM values. We showed the impact of experimental limitations on jet reconstruction and triggering for reconstructing the principal discriminating variable m_{hh} .

Assuming SM top Yukawa y_t , our *baseline analysis* provides 68% CL constraints of $-1.6 < \kappa_\lambda < 7.9$, while the corresponding constraint for the *neural network analysis* is $-0.8 < \kappa_\lambda < 6.6$. This assumes 3000 fb^{-1} of luminosity, background systematics controlled to 0.3%, 1% and 5% for resolved, intermediate and boosted categories respectively, and no further advances in the analysis strategy. Interestingly, we can extract a 68% CL upper

limit on the top Yukawa of $\kappa_t < 1.22$ assuming $\kappa_\lambda = 1$ despite no dedicated optimisation, though this constraint changes rapidly for different κ_λ values. We quantified that current uncertainties on y_t constraints can modify λ_{hhh} limits by $\sim 20\%$. We caution that these constraints are not necessarily directly comparable with other projections in the literature due to different simplifying assumptions.

While our study suggests that the $4b$ channel is challenging for probing λ_{hhh} compared with for example $bb\gamma\gamma$, information from all independent channels are experimentally important in the final statistical combination. Our conclusions sharpen the experimental requirements to improve λ_{hhh} constraints using the $hh \rightarrow 4b$ channel, providing key performance targets for the HL-LHC programme. Given the dominant reducible $2b2j$ background, reducing the flavour mistag rate is critical, which requires both hardware (upgraded trackers improve vertex resolution) and software (e.g. deep learning) solutions. We quantified the impact of different scenarios of background systematic uncertainties from 5% to statistical-only for all three categories. This showed that controlling systematics on multijet background estimates to demanding percent level or better is crucial for improving λ_{hhh} sensitivity. For scenarios where the boost of the Higgs bosons is low, notably around the boundary of expected sensitivity for positive $\kappa_\lambda \sim 5$ values, maintaining low p_T jet triggers will be particularly important. Experimentally, applying and improving this rich confluence of techniques explored will accelerate progress towards constraining λ_{hhh} beyond gains from increased statistics alone.

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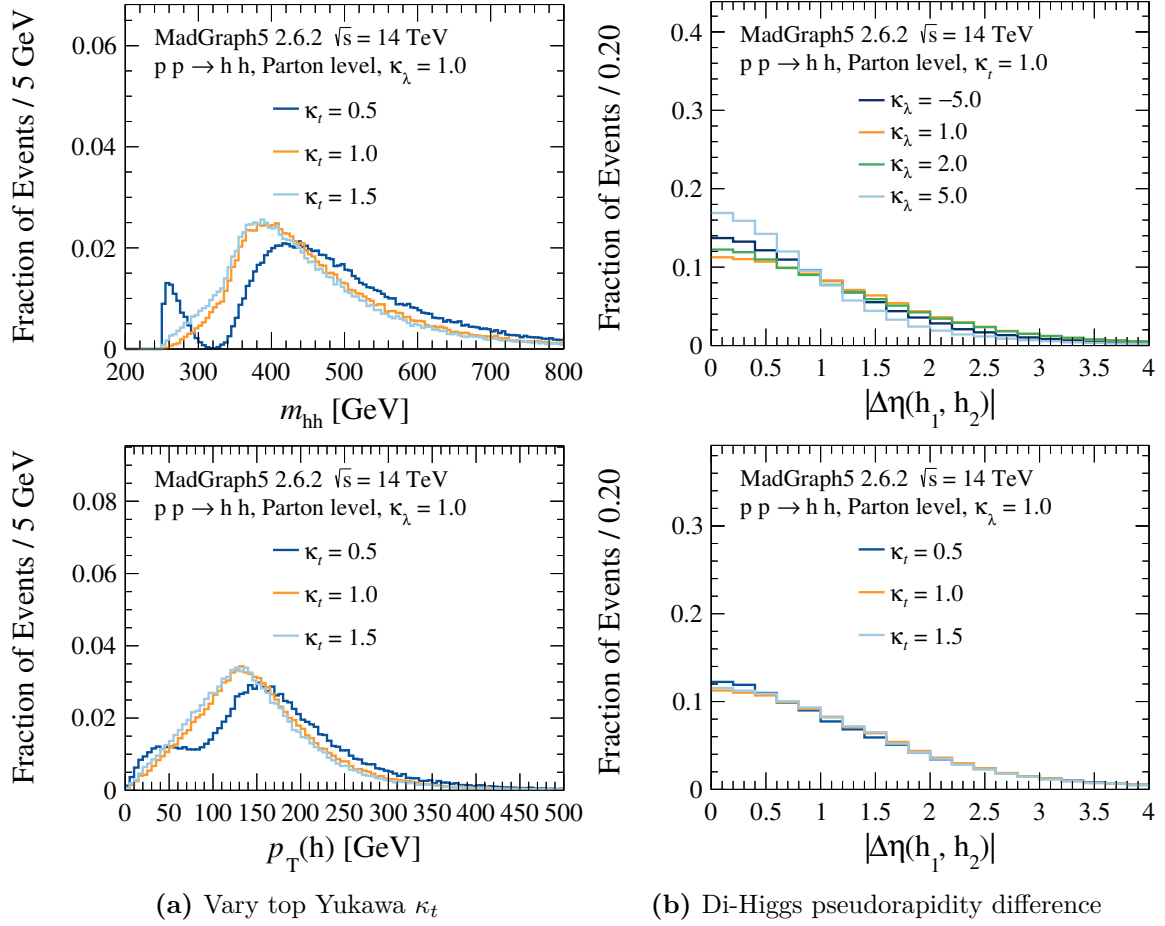


Figure 16. Unit normalised distributions for parton-level signals for different couplings.

A Additional kinematic distributions

This appendix collects additional kinematic distributions. Figure 16a shows the impact at parton level of top Yukawa variations for fixed $\kappa_\lambda = 1$. In particular, $\kappa_t = 1.5$ has little impact on the m_{hh} shape, but lowering κ_t to 0.5 makes the destructive interference induce a similar m_{hh} shape as $(\kappa_\lambda, \kappa_t) = (2, 1)$. Figure 16b displays the rapidity difference between the Higgs bosons $|\Delta\eta(h_1, h_2)|$, which has mild discriminating power between couplings.

Figure 17 displays distributions at reconstructed level for both backgrounds and signals in the *baseline analysis* signal regions normalised to cross-sections. The subleading Higgs candidate $p_T(h_2^{\text{cand}})$ and transverse momentum of the di-Higgs system $p_T(hh)$ are shown. We note that in the resolved category, the $p_T(h_2^{\text{cand}})$ retains some discrimination power between different κ_λ hypotheses. The $p_T(hh)$ is also lower for the resolved than the intermediate and boosted analyses. This suggests that a modest amount of the boost that the Higgs bosons receives arises from the recoil of the di-Higgs system against other activity such as initial state radiation jets.

Figure 18 displays distributions at reconstructed level for both backgrounds and signals normalised to cross-sections. The *baseline analysis* signal region selection is applied

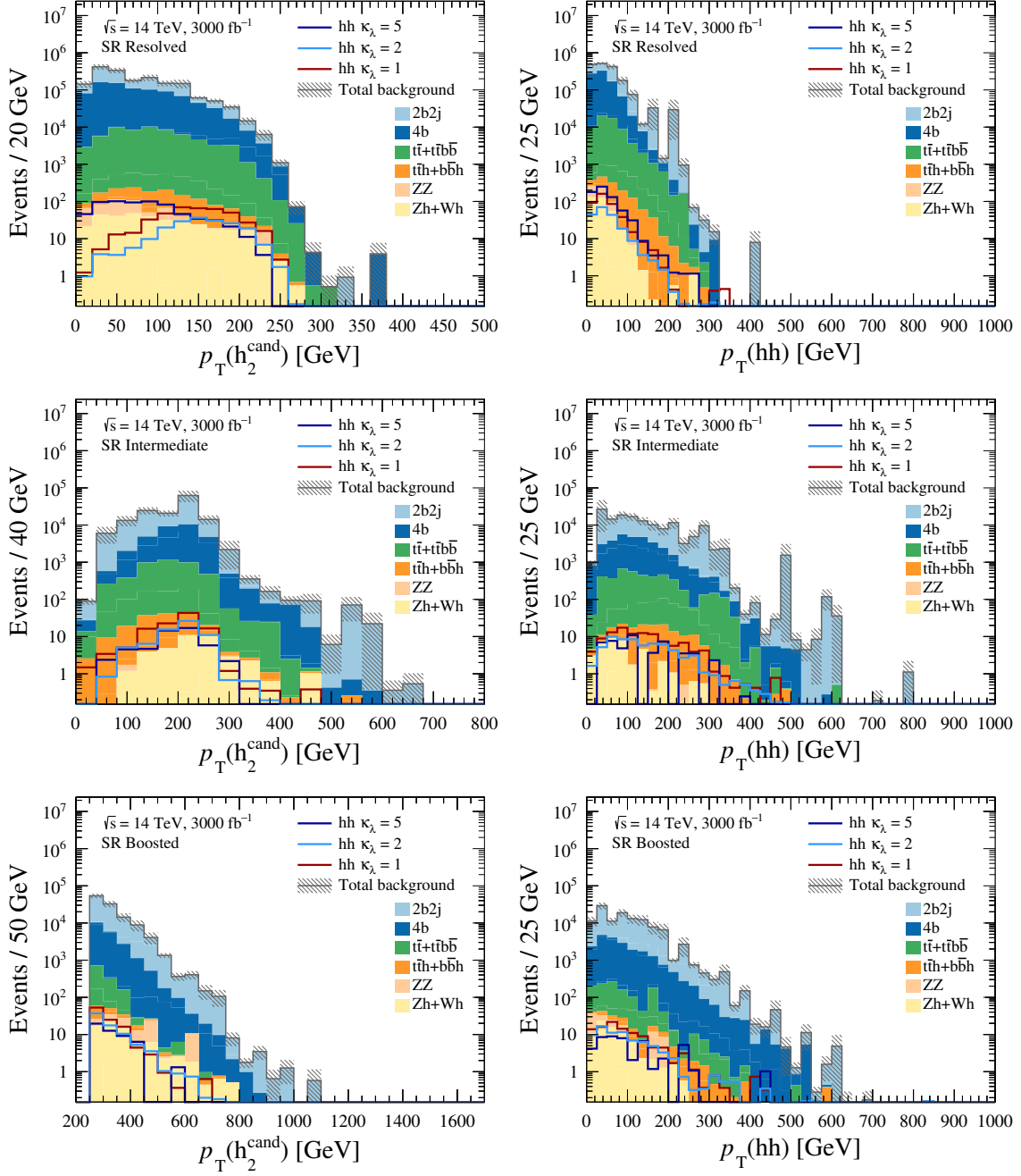


Figure 17. Distributions of signal (lines) and backgrounds (stacked filled) for (left) the transverse momentum of the subleading Higgs candidate $p_T(h_2^{\text{cand}})$ and (right) the di-Higgs system $p_T(hh)$ for the *baseline analysis* in the (upper) resolved, (middle) intermediate and (lower) boosted categories.

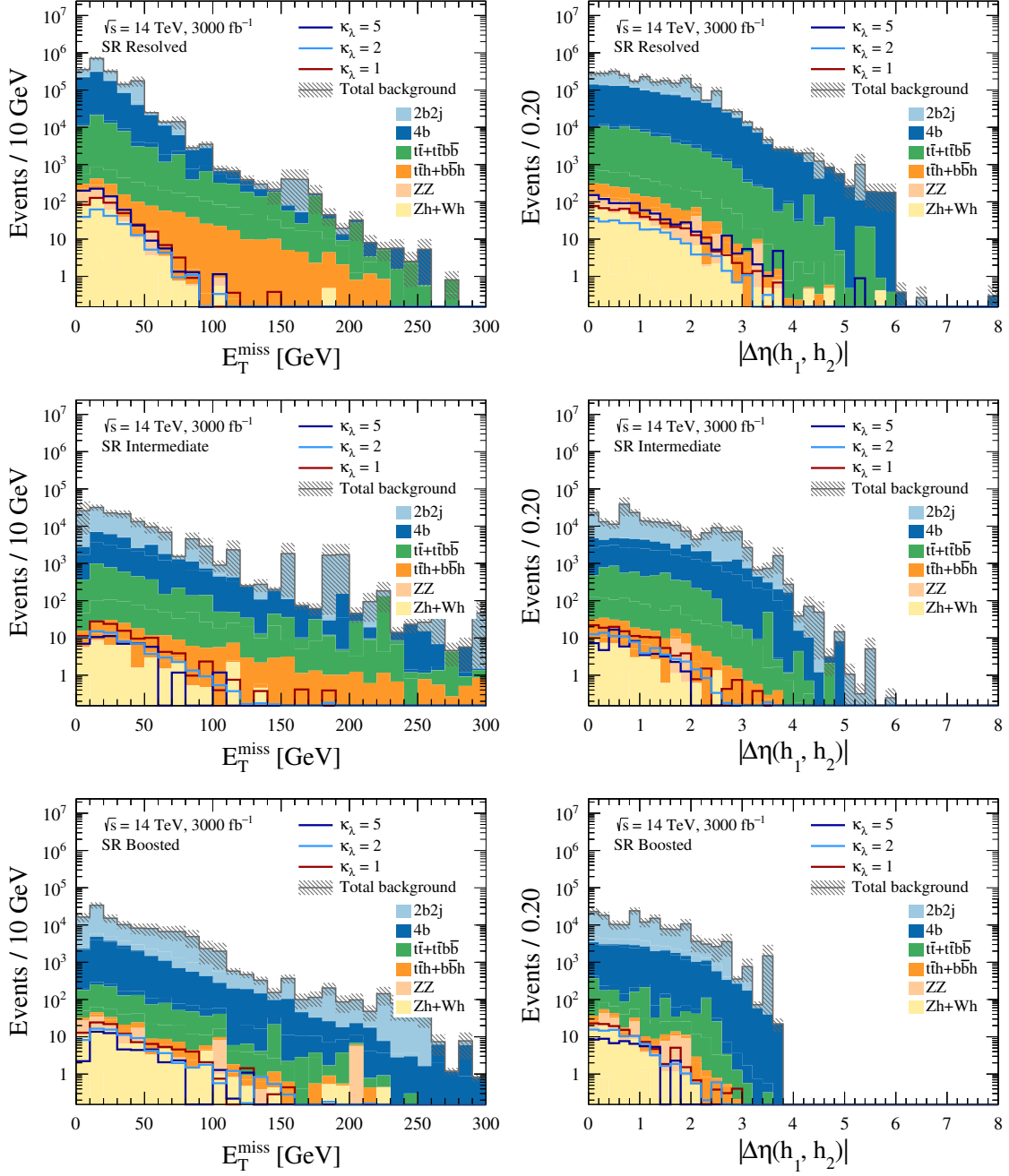


Figure 18. ‘ $N - 1$ ’ distributions with all signal region event selection applied except for the variable being plotted. Displayed are various signal (lines) and backgrounds (stacked filled) for (left) the magnitude of the missing transverse momentum E_T^{miss} and (right) the pseudorapidity difference of the two Higgs candidates $\Delta\eta_{hh}$ for the *baseline analysis* in the (upper) resolved, (middle) intermediate and (lower) boosted categories.

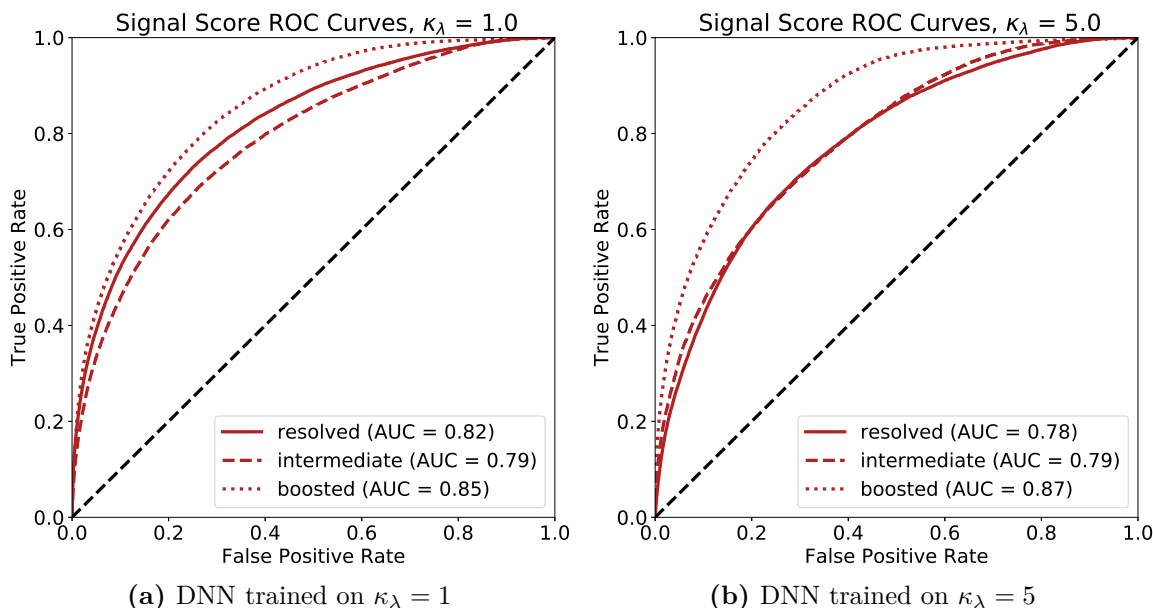


Figure 19. Receiver operator characteristic (ROC) curves for the DNN trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$. These quantify true positive vs false positive classification rates for the signal events in the resolved (solid), intermediate (dashed), and boosted (dotted) categories. Values for the area under the curve (AUC) are displayed in the legend.

except for requirements on the variable being plotted. The magnitude of the missing transverse momentum E_T^{miss} and the pseudorapidity difference of the two Higgs candidates $|\Delta\eta(h_1, h_2)|$ is shown. A mild sensitivity to different κ_λ values is observed for $|\Delta\eta(h_1, h_2)|$, however, this is subdominant compared to the discriminating power of m_{hh} or $p_T(h)$. Other variables employed in our analysis selection are used to reduce background and are physically insensitive to κ_λ . For example, we cut on E_T^{miss} to reduce semileptonic $t\bar{t}$ which has neutrinos in the final state, but this has no physical correlation with kinematics that are affected by modifying κ_λ .

B Neural network performance

This appendix provides further performance benchmarks of the *neural network analysis* for the DNN trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$ signals. All the plots in this section are made with an independent set of events compared with those used for training. The selection criteria applied for the DNN training, and also for all plots in this section, are looser than the final analysis selection in order to retain sufficient event statistics for reliable training. One significant difference is that this dataset does not include any requirements on the number of b -tagged jets.

Figure 19 shows the true positive vs false positive rate known as receiver operator characteristic (ROC) curves. These are computed on the whole test set (i.e. the set of events *not* used for training) for these networks in the resolved, intermediate, and boosted categories for signal classification only. In the general case of multi-class classification, the

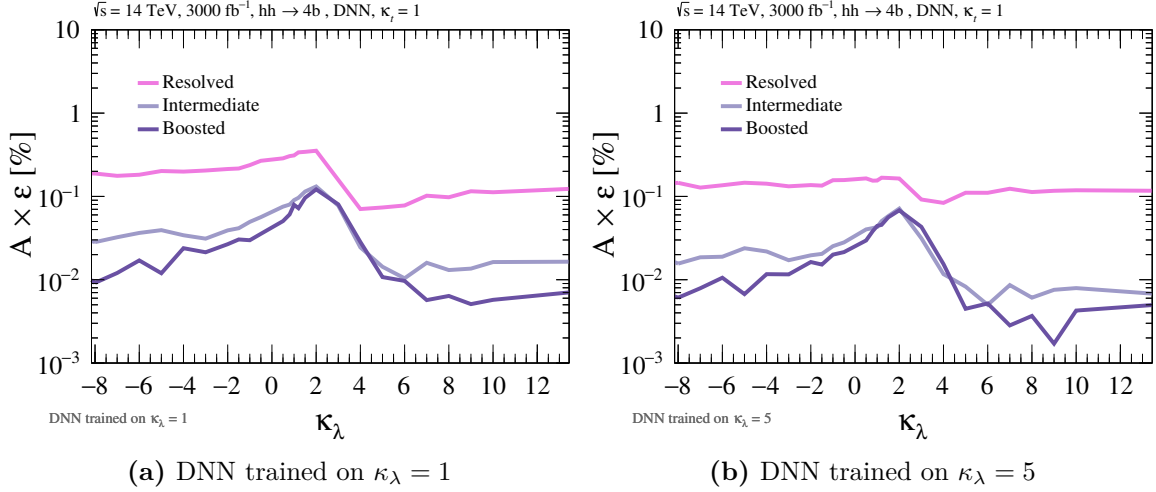


Figure 20. Signal acceptance times efficiency $A \times \varepsilon$ in percent for the *neural network analysis* trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ after the signal score requirement $p_{\text{signal}}^{\text{DNN}} > 0.75$. This is shown for the resolved (pink), intermediate (lilac) and boosted (purple) channels. The $A \times \varepsilon$ is equivalent to the number of signal events S divided by initial number of events $\sigma \times \mathcal{L}$.

true positive rate is defined by the fraction of events where an event from sample X is classified correctly as an event from X divided by the total number of events in X

$$R_{\text{positive}}^{\text{true}} = \frac{N(X \text{ classified as } X)}{N(\text{total } X)}, \quad X \in \{\text{multijet}, t\bar{t}, \text{signal}\}. \quad (\text{B.1})$$

The false positive rate is defined by the fraction of events where an event not from sample X is classified as an event from sample X divided by the total number of events not from sample X

$$R_{\text{positive}}^{\text{false}} = \frac{N(Y \text{ or } Z \text{ classified as } X)}{N(\text{total } Y + Z)}, \quad X, Y, Z \in \{\text{multijet}, t\bar{t}, \text{signal}\}. \quad (\text{B.2})$$

Figures 20 also displays the signal acceptance times efficiency $A \times \varepsilon$ as a function of κ_λ for the different categories in the *neural network analysis*. One can see that the signal acceptance is higher around $\kappa_\lambda \sim 5$ when the DNN is trained on $\kappa_\lambda = 5$.

Finally, figure 21 shows the ranked feature importance for the *neural network analysis* trained on $\kappa_\lambda = 1$. Discussion of this figure follows that presented in section 3.2.1 of the main text.

B.1 Heuristics of Shapley values

Here we provide a heuristic overview of the SHAP method for model interpretation; for technical details and definitions, see ref. [88]. The SHAP method is based on Shapley values ϕ , which were originally introduced in the context of game theory. These calculate the model prediction (neural network score), $f(S \cup x)$, when including a variable x (feature) from a set S . The prediction is computed without this variable $f(S)$. From these, the difference, $f(S \cup x) - f(S)$, intuitively characterises the importance of that variable and

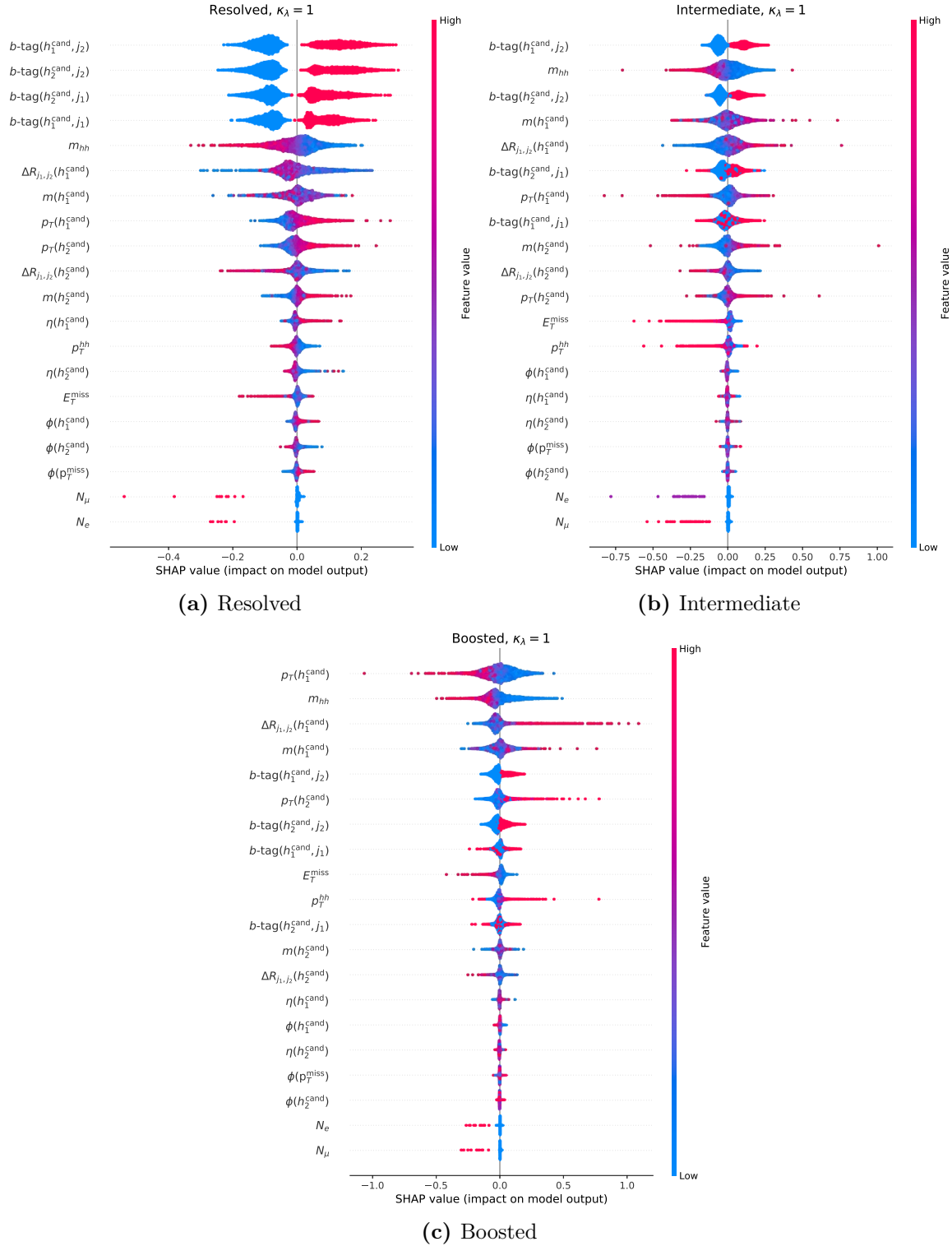


Figure 21. SHAP value plots representing the ranked variable (feature) importance for the DNN trained on $\kappa_\lambda = 1$ signals for the (a) resolved, (b) intermediate and (c) boosted categories. The DNN models are ranked by their average absolute SHAP value, which is plotted on the x -axis. The colour scale indicates the value of the feature on the specific event for which the SHAP value is plotted.

its impact on the model prediction. Shapley variables satisfy other conditions: a) the total value of ϕ is additive in that it equals the sum of the predictions when all the variables are included, b) variables that contribute the same to a prediction have the same to ϕ , c) if the variable does not contribute any change to the prediction f then it adds nothing to ϕ . The predictions are computed for each variable in the set to compare their importance systematically. SHAP (SHapley Additive exPlanation) considers Shapley values of a conditional expectation function of the original model. Formal definitions and prescriptions are found in ref. [88].

B.2 Network correlation plots

We validate the architecture and training of the *neural network analysis* has the expected behaviour by testing their performance on a statistically independent set of signal and background events. Figure 22 show the leading Higgs candidate mass as a function of the signal score trained on $\kappa_\lambda = 1$ and $\kappa_\lambda = 5$ signals when tested on the corresponding $pp \rightarrow hh$ signal that exclusively decays to b -quarks. For the resolved category, a prominent peak appears at high signal score trained on $\kappa_\lambda = 1$ and around $m(h_1^{\text{cand}}) \sim 125$ GeV. There is a tail of events at lower signal score where it is difficult to kinematically distinguish the signal from background.

An interesting feature is observed in figure 22 depicting the $\kappa_\lambda = 5$ DNN score computed in the $\kappa_\lambda = 5$ signal sample versus the leading Higgs candidate mass. Three distinct peaks can be seen. We find each peak corresponds to events with different numbers of b -tags. The most prominent peak with the highest signal score corresponds to events with three or four b -tags. Meanwhile, the two peaks with signal scores around 0.3 and 0.5 correspond to events with one and two b -tags respectively.

Figures 23 show similar correlation plots, now testing on a $t\bar{t}$ sample. We see low values of signal scores, indicating that the networks are effective at classifying this background. More structure is seen in these two-dimensional plots due to the inclusive decay of the b -quarks in this sample. Several of these plots feature two distinct peaks, corresponding to semi-leptonic and hadronic b -decays. Hadronic decays correspond to the peak with higher signal score, since these are more similar to the $hh \rightarrow 4b$ signal. Furthermore, the plots for the intermediate and boosted categories feature two distinct peaks in the leading Higgs mass around ~ 80 GeV and ~ 170 GeV. This suggests that the large jet is capturing the boosted decay products of the W boson and top quarks.

Figures 24 and 25 show the corresponding correlation plots for the di-Higgs invariant mass m_{hh} as a function of DNN signal score. We see that the signal score is not strongly correlated with m_{hh} .

C Additional χ^2 distributions

This appendix collects χ^2 distributions supplementing those in the main text.

Figure 26 shows the discrimination power χ_{ij}^2 matrix between different κ_λ hypotheses for $\kappa_t = 1$ for the *baseline analysis* and *neural network analysis* trained on $\kappa_\lambda = 1$. We find a similar pattern of regions with lower χ_{ij}^2 values as figure 13 in the main text, where further discussion can be found.

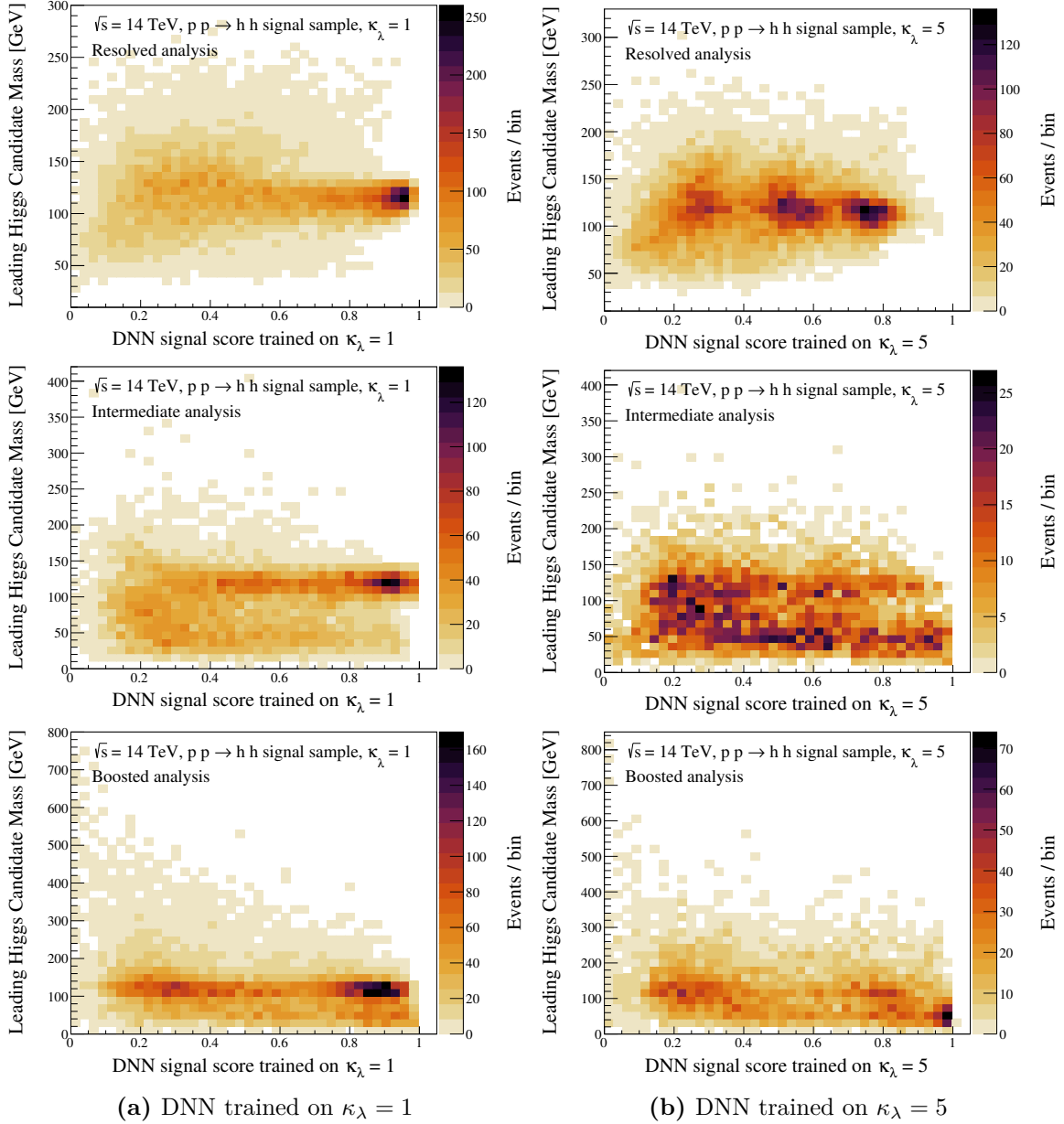


Figure 22. The leading Higgs candidate mass vs neural network scores trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal sample for the (upper) resolved, (middle) intermediate, and (lower) boosted analyses. The test samples used to make these distributions are an independent set of (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal events.

Figure 27 shows a summary of the 68% CL limits for different systematics when the DNN is trained on the $\kappa_\lambda = 1$ signal. The qualitative features are similar to the corresponding figure 12 where the DNN is trained on the $\kappa_\lambda = 5$ signal. As noted in the main text, the DNN trained on $\kappa_\lambda = 1$ for the intermediate and boosted categories are slightly more constraining than that trained on $\kappa_\lambda = 5$.

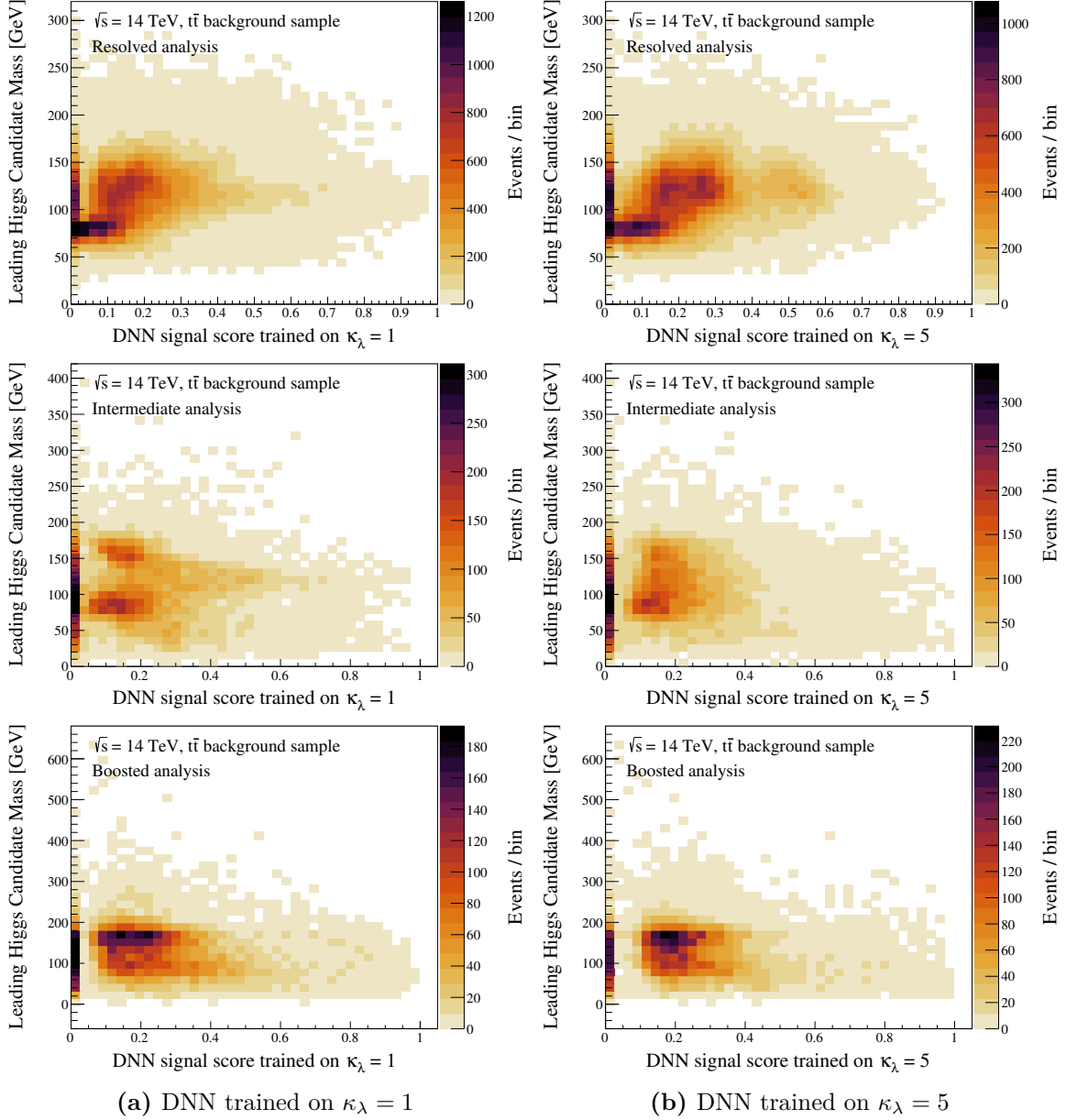


Figure 23. The leading Higgs candidate mass vs neural network scores trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal sample for the (upper) resolved, (middle) intermediate, and (lower) boosted analyses. The test samples used to make these distributions are an independent set of $t\bar{t}$ events.

Figure 28 shows the χ^2 distributions for the *baseline* and *neural network* analyses trained on $\kappa_\lambda = 5$ separated by resolved, intermediate and boosted categories. They assume 3000 fb^{-1} of luminosity and 0.3%, 1% and 5% systematic uncertainties for the resolved, intermediate and boosted categories, respectively. The statistical combination of the three categories results in the χ^2 distributions displayed in figure 14 of the main text, where further discussion is found.

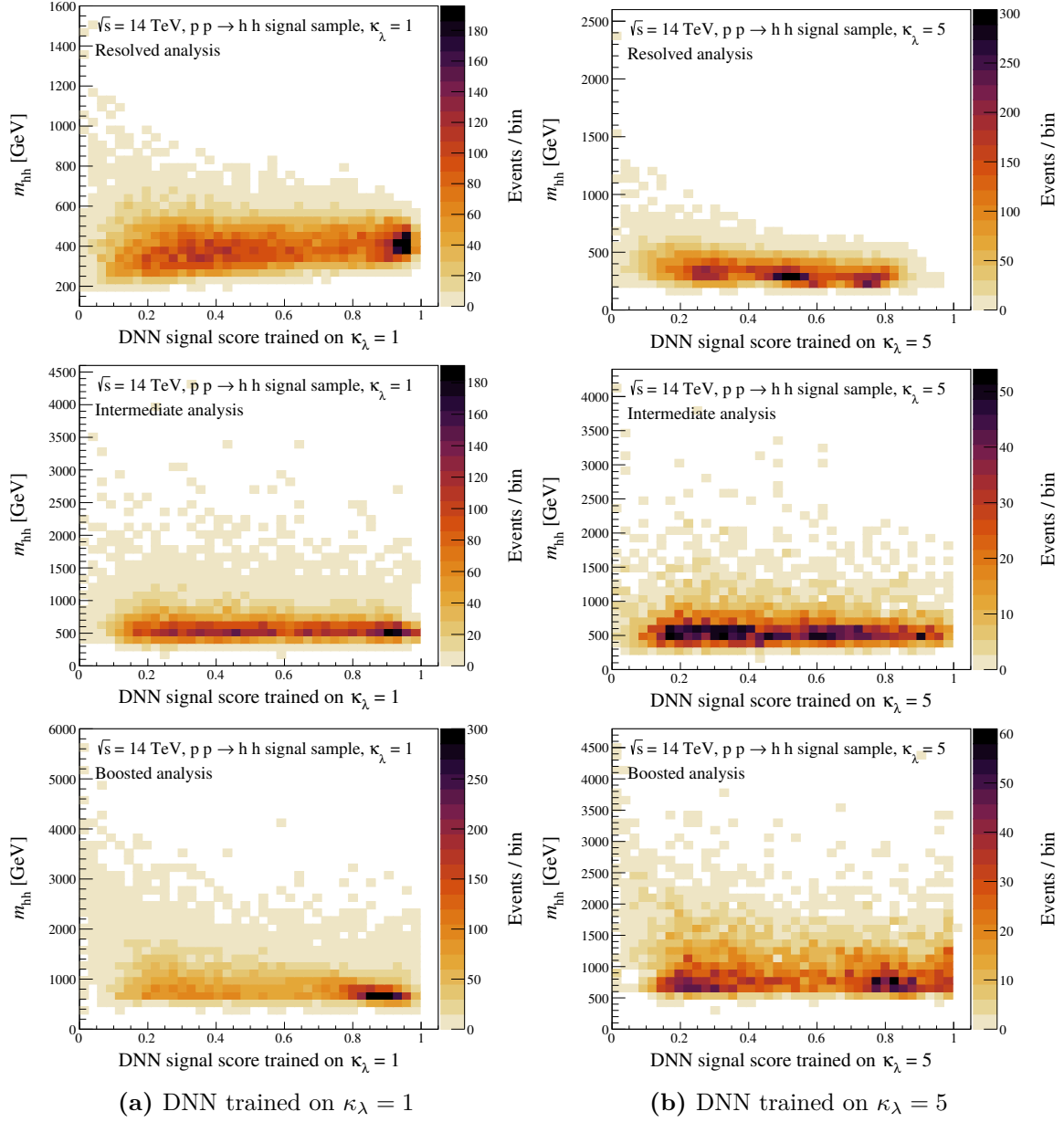


Figure 24. The di-Higgs invariant mass m_{hh} vs neural network scores trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal sample for the (upper) resolved, (middle) intermediate, and (lower) boosted analyses. The test samples used to make these distributions are an independent set of (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal events.

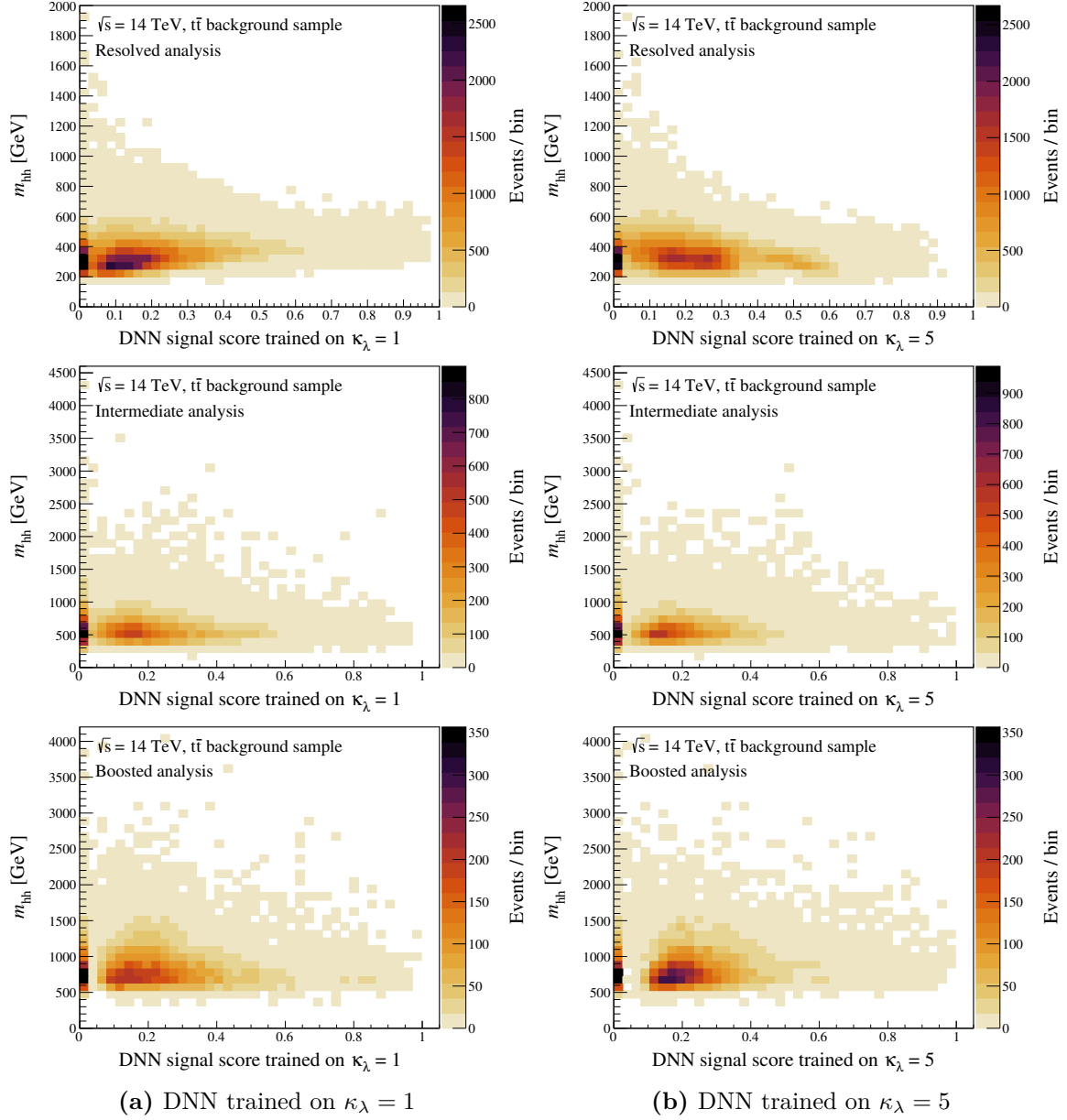


Figure 25. The di-Higgs invariant mass m_{hh} vs neural network scores trained on (a) $\kappa_\lambda = 1$ and (b) $\kappa_\lambda = 5$ signal sample for the (upper) resolved, (middle) intermediate, and (lower) boosted analyses. The test samples used to make these distributions are an independent set of $t\bar{t}$ events.

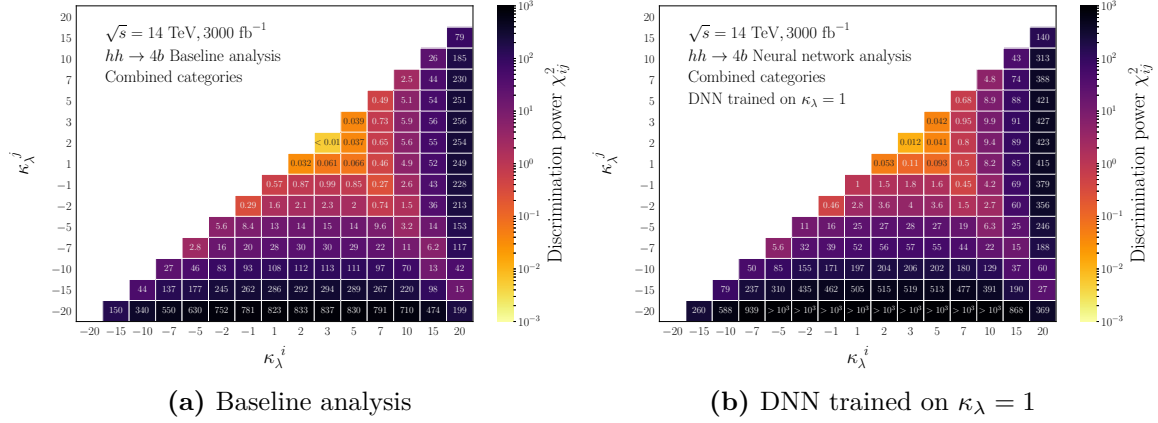


Figure 26. Discrimination power χ_{ij}^2 (higher is better) between two self-coupling values κ_λ^i vs κ_λ^j . This is displayed for the (a) *baseline* and (b) *neural network* trained on $\kappa_\lambda = 1$ analyses. The requirement on the signal score $p_{\text{signal}} > 0.75$ is imposed for the DNN. The χ_{ij}^2 is calculated from eq. (3.1). This assumes $\kappa_t = 1.0$ and background systematics of $\zeta_b = 0.3\%$, 1% and 5% for the resolved, intermediate and boosted categories, respectively.

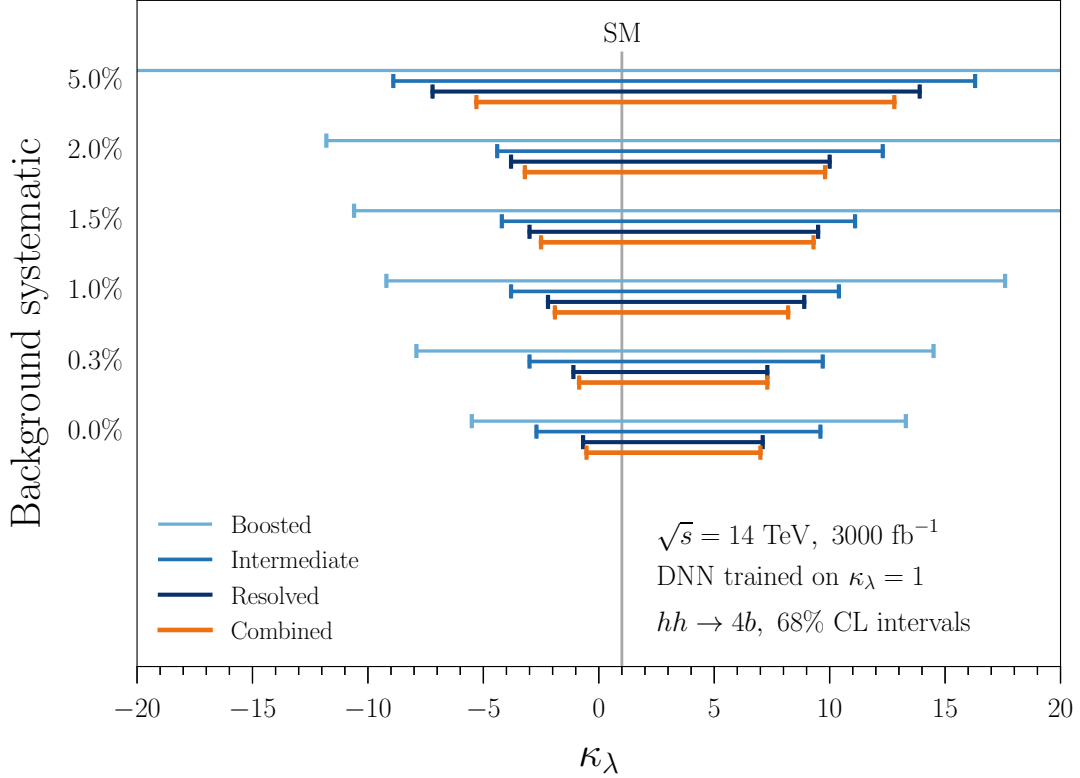


Figure 27. Summary of 68% CL intervals ($\chi^2 < 1$) on κ_λ with fixed $\kappa_t = 1$ for different assumed background systematics from 0% to 5% for the *neural network analysis* trained on the $\kappa_\lambda = 1$ signal. This is shown for the resolved (dark blue), intermediate (medium blue), boosted (light blue) categories along with their combination (orange). The luminosity is assumed to be 3000 fb $^{-1}$. Lines without endcaps mean the 68% CL limit is outside the considered range of $\kappa_\lambda \in [-20, 20]$. The vertical grey line denotes the SM value of κ_λ . The nominal systematics assumed are shown in table 6 of the main text.

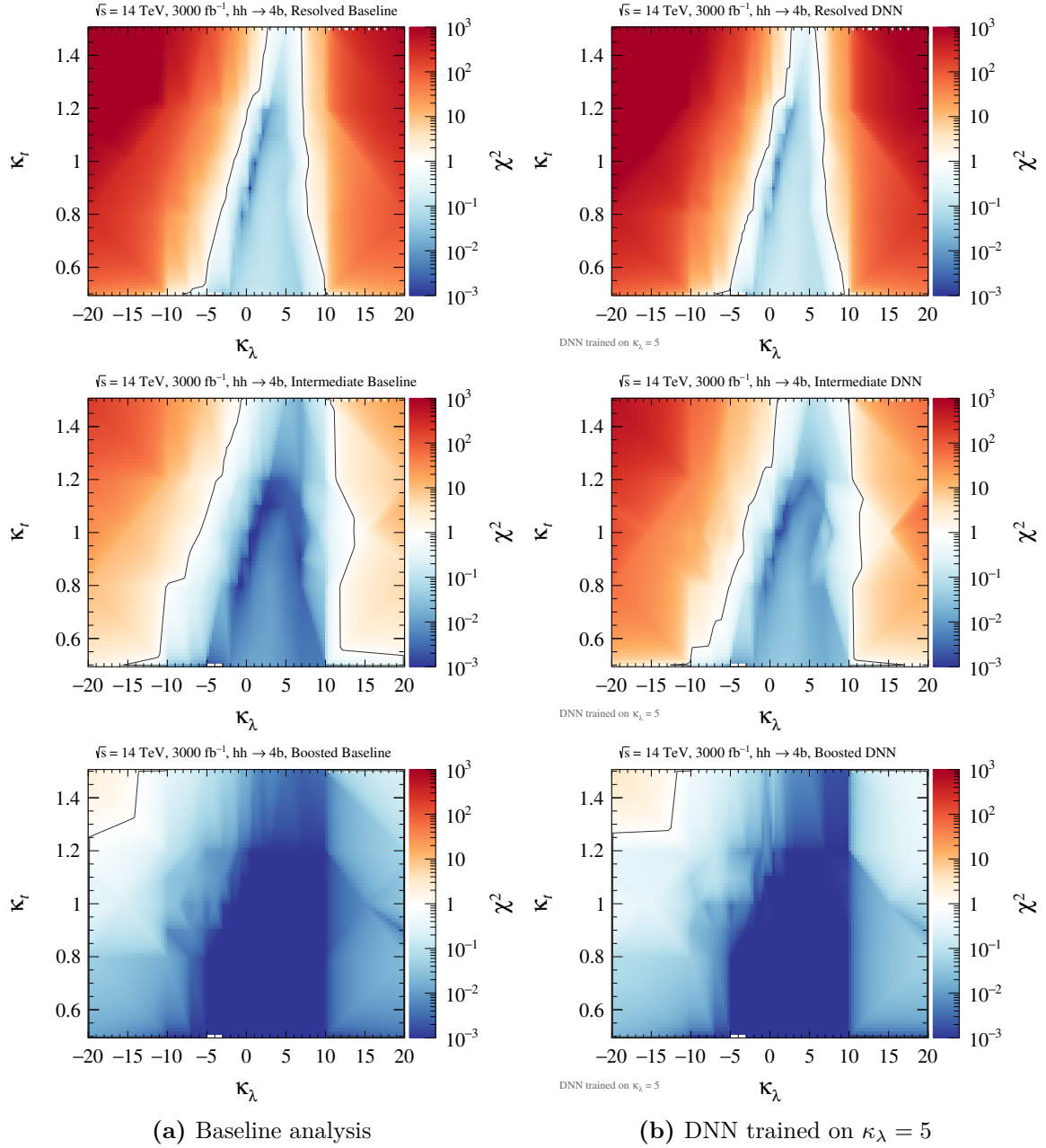


Figure 28. The χ^2 distributions for the (a) *baseline analysis* and (b) *neural network analysis* trained on the $\kappa_\lambda = 5$ signal. This is displayed in the κ_t vs κ_λ plane at $\mathcal{L} = 3000 \text{ fb}^{-1}$. This is shown for the resolved (upper), intermediate (middle), boosted (lower) categories assuming 0.3%, 1% and 5% systematic uncertainties, respectively. The grey contour indicates $\chi^2 = 1$ corresponding to 68% CL.

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